

THE FRENET FORMULAE OF THE RIEMANN-OTSUKI
SPACE

Djerdji F. Nadj

*University of Sopron, Department of
Mathematics, H-9401 Sopron, Hungary*

ABSTRACT

In this paper the Frenet formulae of the Riemann-Otsuki space with respect to covariant and contravariant part of the connection are obtained.

INTRODUCTION

The basis of the theory of Otsuki spaces has been laid down by T. Otsuki and A. Moór. The metric used determined that the observed space is of Weyl-Otsuki's or of Riemann-Otsuki's kind. In this paper we shall consider the Riemann-Otsuki space and we shall determine the Frenet formulae with respect to the co-resp. contravariant part of the connection. According to the following observation, we get that only in the contravariant part of the connection the Frenet formulae of the $R-O_n$ space are different from the known Frenet formulae of Riemannian geometry. The difference came from the fact that in Otsuki spaces $D\delta^i_j \neq 0$ holds.

AMS Mathematics Subject Classifications (1980): Primary 53B05, Secondary 53B15.

Key words and phrases: Riemann-Otsuki spaces, Frenet formulae.

In all Otsuki spaces we have, with respect to the local coordinates x^i of an n -dimensional differentiable manifold, an *a-priori* given tensor P_j^i such that $\det\|P_j^i\| \neq 0$ holds and the inverse tensor Q_j^i exists so that $P_j^i Q_r^j = \delta_r^i$. In the metric Otsuki spaces the metric tensor g_{ij} ($\det\|g_{ij}\| \neq 0$) is given so that in the $W - O_n$ Weyl-Otsuki space $\nabla_k g_{ij} = \gamma_k g_{ij}$, but in the $R - O_n$ (Riemann-Otsuki space)

$$(0.1) \quad \nabla_k g_{ij} = 0$$

holds. In Otsuki spaces the covariant differential of the tensor T_j^i is defined by

$$(0.2) \quad DT_j^i = P_a^i P_j^b DT_b^a = P_a^i P_j^b (\partial_k T_b^a + \overset{a}{\Gamma}_{r k} T_b^r - \overset{r}{\Gamma}_{b k} T_r^a) dx^k.$$

The Leibnitz formula *does not hold* for this differential. The differential \bar{D} is the *basic* covariant differential. The different coefficients of the connection are characteristic of the Otsuki spaces, and here are

$$(0.3) \quad \delta_{j|k}^i = \overset{i}{\Gamma}_{j k}^i - \overset{i}{\Gamma}_{j k}^i \neq 0.$$

The coefficient of the connection $\overset{i}{\Gamma}_{j k}^i$ was determined from the relation (0.1) and the coefficients of connection $\overset{i}{\Gamma}_{j k}^i$ are got from

$$(0.4) \quad \partial_k P_j^i + \overset{i}{\Gamma}_{a k}^i P_j^a - P_a^i \overset{a}{\Gamma}_{j k}^a = 0.$$

This relation is known as *Otsuki's relation*.

In Otsuki spaces it is possible to determine the covariant differentials D and \bar{D} with respect only to the co-resp. contravariant part of the connection. So

$$(0.5) \quad \bar{DT}_j^i = \nabla_k T_j^i dx^k = (\partial_k T_j^i + \overset{i}{\Gamma}_{r k}^i T_j^r - \overset{r}{\Gamma}_{j k}^r T_r^i) dx^k$$

holds. For this basic covariant differential the Leibnitz

formula holds. The basic covariant differential $\overset{\circ}{\bar{D}}$ can be defined in the same way. It is characteristic that the basic covariant differential $\overset{\circ}{\bar{D}}$ is identical in the case of contravariant indices with the basic covariant differential \bar{D} , and similarly in the case of covariant indices the basic covariant differential $\overset{\circ}{\bar{D}}$ is identical with the basic covariant differential \bar{D} .

In the following we shall use the relations

$$(0.6) \quad \overset{\circ}{\bar{D}}g_{ij} = dg_{ij} - (\overset{\circ}{\Gamma}_{ik}^r g_{rj} + \overset{\circ}{\Gamma}_{jk}^r g_{ir}) dx^k,$$

$$(0.7) \quad \overset{\circ}{\bar{D}}g_{ij} = dg_{ij} - (\overset{\circ}{\Gamma}_{ik}^r g_{rj} + \overset{\circ}{\Gamma}_{jk}^r g_{ir}) dx^k = 0,$$

$$(0.8) \quad \overset{\circ}{\bar{D}}g^{ra} = -g^{ia} g^{jr} (\overset{\circ}{\bar{D}}g_{ij}),$$

$$(0.9) \quad \overset{\circ}{\bar{D}}g^{ra} = 0.$$

1. THE FRENET FORMULAE WITH RESPECT TO THE CONTRAVARIANT COMPONENTS OF THE VECTORS

Let the point P of the curve $C : x^i = x^i(s)$ be given where s is the arch length parameter. In that point $y^i := \frac{dx^i}{ds}$ are the components of the unit tangent vector y . Applying the basic covariant differential $\overset{\circ}{\bar{D}}$ on the relation

$$(1.1) \quad g_{ij} y^i y^j = 1$$

using the Leibnitz formula and the symmetry of the tensor g_{ij} we get

$$(1.2) \quad (\overset{\circ}{\bar{D}}g_{ij}) y^i y^j + 2g_{ij} y^i (\overset{\circ}{\bar{D}}y^j) = 0.$$

From relations (0.6) and (0.7), it follows that $\overset{\circ}{\bar{D}}g_{ij} = -2g_{r(i} \overset{\circ}{\bar{D}}\delta_{j)}^r$. Substituting it in relation (1.2) and using that $\overset{\circ}{\bar{D}}y^j = \overset{\circ}{\bar{D}}v^j$, we get

$$g_{ij} v_0^i (\bar{D}v_0^j - v_0^r \bar{D}\delta_r^j) = 0.$$

Let v_1^i be the components of the unit vector v_1 so that

$$(1.3) \quad a) \quad v_1^j := \frac{1}{\kappa(s)} (\bar{D}v_0^j - v_0^r \bar{D}\delta_r^j), \kappa_1 > 0; \quad b) \quad g_{ij} v_1^i v_1^j = 1$$

holds. From the above relations we can see, that $v_1 \perp v_0$ and

$$(1.4) \quad \kappa_1(s) = \left(g_{rq} (\bar{D}v_0^r - v_0^a \bar{D}\delta_a^r) (\bar{D}v_0^q - v_0^b \bar{D}\delta_b^q) \right)^{\frac{1}{2}}$$

and

$$(1.5) \quad \bar{D}v_0^j = \kappa_1 v_1^j + v_0^q \bar{D}\delta_q^j$$

hold. (1.5) is the *first Frenet formula* of the basic covariant differential applied on the *contravariant* components of the vectors.

Applying now the basic covariant differential \bar{D} on relation (1.3.b) with a calculation like the one above we get that $g_{ij} v_1^i (\bar{D}v_1^j - v_1^r \bar{D}\delta_r^j) = 0$. This means that the vector v_1 is orthogonal to the direction of $\bar{D}v_1^j - v_1^r \bar{D}\delta_r^j$, and if v_2^i denotes the components of the unit vector orthogonal on the plain determined by the vectors v_0 and v_1 , it follows that

$$(1.6) \quad \bar{D}v_1^j - v_1^r \bar{D}\delta_r^j = \alpha v_0^j + \kappa_2 v_2^j$$

holds. Now we shall use that $v_0 \perp v_2$ and apply the covariant differential \bar{D} on the relation $g_{ij} v_0^i v_2^j = 0$. A calculation like the one above with the substitution $\bar{D}v_0^j$ from (1.5) and $\bar{D}v_1^j$ from (1.6) according to (1.1) and (1.3 b) gives

$$(1.7) \quad \alpha = -\kappa_2(s).$$

Substituting it in (1.6), we get

$$(1.8) \quad \bar{D}v_1^j = -\kappa_{10}^j v_1^j + \kappa_{22}^j v_1^j + v_1^r \bar{D}\delta_r^j.$$

From this relation there is $v_2^i = \frac{1}{\kappa_2(s)} (\bar{D}v_1^i + \kappa_{10}^i v_1^i - v_1^r \bar{D}\delta_r^i)$ and

$$\kappa_2(s) = \left(g_{ij} (\bar{D}v_1^j + \kappa_{10}^j v_1^j - v_1^r \bar{D}\delta_r^j) (\bar{D}v_1^i + \kappa_{10}^i v_1^i - v_1^r \bar{D}\delta_r^i) \right)^{\frac{1}{2}}.$$

Now we can formulate

Lemma 1. If v is the unit vector, then $(\bar{D}v^j - v^r \bar{D}\delta_r^j) g_{ij} v^i = 0$ and $(\bar{D}v_j - v_r \bar{D}\delta_j^r) g^{ij} v_i = 0$ holds.

(The proof of the covariant case follows in §2.)

We shall now make the following generalization.

Let for mutually orthogonal unit vectors v_ℓ ($\ell = 0, \dots, p-1$) be

$$(1.9) \quad \bar{D}v_\ell^j = -\kappa_{\ell-1}^j v_{\ell-1}^j + \kappa_{\ell+1}^j v_{\ell+1}^j + v_\ell^r \bar{D}\delta_r^j$$

so that $\kappa_0 = 0$, and if $q = 1, \dots, p-1$ then

$$(1.10) \quad \kappa_q = \left(g_{ij} (\bar{D}v_{q-1}^j + \kappa_{q-1}^j v_{q-2}^j - v_{q-1}^r \bar{D}\delta_r^j) (\bar{D}v_{q-1}^i + \kappa_{q-1}^i v_{q-2}^i - v_{q-1}^t \bar{D}\delta_t^i) \right)^{\frac{1}{2}}$$

holds. We construct the vector v_p^i ($p < n$) so that

$$(1.11) \quad v_p^i := \frac{1}{\kappa_p} (\bar{D}v_{p-1}^i + \kappa_{p-2}^i v_{p-2}^i - v_{p-1}^r \bar{D}\delta_r^i).$$

According to the above Lemma $g_{ij} v_p^i (\bar{D}v_p^j - v_p^r \bar{D}\delta_r^j) = 0$ holds and we can write the linear combination

$$(1.12) \quad \bar{D}v_p^j - v_p^r \bar{D}\delta_r^j = \alpha_0 v_0^j + \alpha_1 v_1^j + \dots + \alpha_{p-1} v_{p-1}^j + \alpha_{p+1}^j v_{p+1}^j.$$

Contracting this with $g_{ij} v_\ell^i$, using that $v_\ell \perp v_m$ ($m \neq \ell$) and v_ℓ is the unit vector, we get

$$\alpha_\ell = g_{ij} v_\ell^i (\bar{D}_p v^j - v_p^r \bar{D}_r v^j).$$

With respect to (1.9) $\alpha_\ell = -g_{ij} v_{\ell+1}^i v_{\ell+1}^j \kappa_{\ell+1}$, i.e. if $\ell \neq p-1$, then $\alpha_\ell = 0$, and if $\ell = p-1$, then $\alpha = -\kappa_p$ holds. Substituting this in (1.12), it follows

$$(1.13) \quad \bar{D}_p v^j = -\kappa_p v_{p-1}^j + \kappa_{p+1} v_{p+1}^j + v_p^a \bar{D}_a v^j$$

and we can formulate

Theorem 1. *If $C : x^i(s)$ is the curve of an $R - O_n$ space and $v_\ell, \ell = 0, \dots, p-1$, ($p < n$) are mutually orthogonal unit vectors which satisfy the relation (1.9) and v_{p+1} is the unit vector orthogonal to all before and $\kappa_0 = 0$, $\kappa_n = 0$ holds, then the vector v_p satisfies the relation (1.9), too.*

If we use Otsuki's covariant differential D , then from the connection $Dv^j = P_a^j \bar{D}_a v^a$ it follows that $\bar{D}_v^a = Q_i^a Dv^i$. Applying this on (1.9), we get

$$(1.14) \quad Dv_\ell^j = P_i^j (-\kappa_{\ell\ell-1} v_{\ell-1}^i + \kappa_{\ell+1\ell+1} v_{\ell+1}^i) + v_\ell^a Q_a^b Dv_b^j$$

with respect to $\ell = 0, \dots, (n-1)$; $\kappa_0 = 0$; $\kappa_n = 0$. We can now state.

Theorem 2. *If in the point M of the curve C in the $R - O_n$ space the mutually orthogonal unit vectors v_0, v_1, \dots, v_{n-1} satisfying relations (1.9) and (1.10) so that $\kappa_0 = 0$ and $\kappa_n = 0$ hold, then (1.14) is the Frenet formula of the curve C of the $R - O_n$ space. This formula is applied with respect to the covariant differential D on the contravariant components of the observed vectors.*

Remark 1. The relation (1.14) is the Frenet formula with respect to the covariant differential \bar{D} , too, applied on the contravariant components of the vectors.

One can see that the difference between (1.14) and the known formula of Riemannian geometry is the covariant differential of the Kronecker- δ . In such a case but not only such special ones where this differential is zero, the formula (1.14) reduces to the known Frenet formula multiplied with the tensor P_j^i .

Now we shall apply the basic covariant differential \bar{D} on (1.1). Using the Leibnitz formula, the symmetry of the tensor g_{ij} and (0.7), we get that $g_{ij} v_0^i (\bar{D}_0^j) = 0$, i.e. $v_0 \perp \bar{D}_0 v_0$ holds. This means that we can construct the unit vector $\bar{v}_1 \perp v_0$ so that

$$(1.15) \quad \bar{D}_0^j = \kappa_1^* \bar{v}_1^j,$$

where the scalar $\kappa_1^*(s)$ satisfies

$$(1.16) \quad \kappa_1^*(s) = (g_{ij} \bar{D}_0^i \bar{D}_0^j)^{\frac{1}{2}} > 0.$$

The relation (1.15) is the first Frenet formula of the basic covariant differential \bar{D} applied on the contravariant components of the vectors. Now we shall prove

Theorem 3. From the connection between the basic covariant differentials \bar{D} and \bar{D} it follows that $v_1 = \bar{v}_1$ and the value of κ is equivalent to the value of κ_1^* .

Proof. It is easy to see that

$$(1.17) \quad \bar{D}_0^i = \bar{D}_0^i - v_0^q \bar{D}_q^i$$

holds. Substituting this in (1.16) we get that according to (1.5) $\kappa_1(s) = \kappa_1^*(s)$ holds. In the following * by $\kappa_1(s)$ we shall

denote that the curvature will be expressed with the aid of the basic covariant differential $\overline{\mathcal{D}}$.

According to the characteristics of the covariant differential $\overline{\mathcal{D}}$, we can state

Theorem 4. *With respect to the basic covariant differential $\overline{\mathcal{D}}$ the Frenet formula of the curve C of the $R - O_n$ space is not different from the known formula of the Riemannian space. If v_0, v_1, \dots, v_{n-1} are in point P of curve C in a suitable way constructed mutually orthogonal unit vectors, then*

$$(1.18) \quad \overline{\mathcal{D}}_l^j v_i^j = P_i^j \left(-\kappa_l^* v_{l-1}^i + \kappa_{l+1}^* v_{l+1}^i \right)$$

is the Frenet formula with respect to the covariant differential applied on the contravariant components of the observed vectors.

2. THE FRENET FORMULAE WITH RESPECT TO THE COVARIANT COMPONENTS OF THE VECTORS

According to the definition $v_{oi} = g_{ij} \frac{dx^j}{ds}$ holds and v_{oi} are the covariant components of the unit tangent vector v_o . Applying the basic covariant differential $\overline{\mathcal{D}}$ on the relation $g^{ij} v_{oi} v_{oj} = 1$, using the Leibnitz formula and relations (0.8), (0.6) and (0.7), we get $g^{ij} v_{oi} (\overline{\mathcal{D}} v_{oj} + v_{ob} \overline{\mathcal{D}} \delta_j^b) = 0$. This proved the second part of Lemma 1 from the first paragraph. Now again, as in the first paragraph, we can construct the mutually orthogonal unit vectors v_l ($l = 0, 1, \dots, n-1$), so that

$$(2.1) \quad \overline{\mathcal{D}} v_l^j = -\kappa_l^{**} v_{l-1}^j + \kappa_{l+1}^{**} v_{l+1}^j - \overline{v}_l^r \overline{\mathcal{D}} \delta_j^r$$

holds with the remarks $\kappa_0^{**} = 0$, $\kappa_n^{**} = 0$, and if $l = 0, \dots, n-2$ then

$$(2.2) \quad \kappa_{l+1}^{**} = \left(g^{ij} \left(\overline{\mathcal{D}} v_l^j + \kappa_l^{**} v_{l-1}^j + \overline{v}_l^r \overline{\mathcal{D}} \delta_j^r \right) \left(\overline{\mathcal{D}} v_{l+1}^i + \kappa_{l+1}^{**} v_{l+1}^i + \overline{v}_{l+1}^q \overline{\mathcal{D}} \delta_i^q \right) \right)^{\frac{1}{2}}$$

We can now formulate

Theorem 5. From the relation $\bar{v}_i = g_{ij} \bar{v}^j$ it follows that the value of κ_{ℓ}^{**} is equal to the value of κ_{ℓ} and $\bar{v}_{\ell} \equiv v_{\ell}$ holds.

Proof. If $\ell = 0$ then according to $\bar{v}_{0i} \equiv v_{0i}$ and $v_{0i} = g_{ij} v_0^j$ using the Leibnitz formula from (2.2), we get

$$\kappa_1^{**} = \left(g^{ij} ((\bar{D}g_{ja})_0^a + g_{ja} (\bar{D}v_0^a) + g_{ra0} v_0^a \bar{D}\delta_j^r) \right. \\ \left. ((\bar{D}g_{ib})_0^b + g_{ib} (\bar{D}v_0^b) + g_{sbo} v_0^b \bar{D}\delta_i^s) \right)^{\frac{1}{2}}.$$

Since $\bar{D}g_{ja} = -2g_s(j\bar{D}\delta_a^s)$ and (1.5) holds according to (1.3 b), it follows that $\kappa_1^{**} = \kappa_1$. Now we shall suppose that $\kappa_{\ell}^{**} = \kappa_{\ell}$ holds, and using the calculation as above, we get that

$$\kappa_{\ell+1}^{**} = (g_{ij} \bar{v}_{\ell+1}^i \bar{v}_{\ell+1}^j \kappa_{\ell+1}^2)^{\frac{1}{2}}$$

Since $\bar{v}_{\ell+1}^i$ are the unit vectors here, the statement of the first part of the theorem follows.

To prove the second part we shall use that $v_0 \equiv \bar{v}_0$. Now we shall suppose that $\bar{v}_p \equiv v_p$ ($p = 0, \dots, \ell$), and from (2.1), using the first part of the theorem, it follows that

$$\bar{v}_{\ell+1}^i = \frac{1}{\kappa_{\ell+1}} (\bar{D}v_{\ell+1}^i + \kappa_{\ell} \frac{v_{\ell+1}^i}{v_{\ell+1}^j} + v_{\ell+1}^a \bar{D}\delta_j^a).$$

Since $v_i = g_{ij} v^j$, using the Leibnitz formula and the calculation as above, contracting by g_{ia} , we get that

$$\bar{v}_{\ell+1}^r = \frac{1}{\kappa_{\ell+1}} (\bar{D}v_{\ell+1}^r + \kappa_{\ell} \frac{v_{\ell+1}^r}{v_{\ell+1}^a} - v_{\ell+1}^a \bar{D}\delta_a^r),$$

i.e. according to (1.11) $\bar{v}_{\ell+1}^r = v_{\ell+1}^r$ holds.

In the following κ_{ℓ}^{**} by κ_{ℓ} we shall denote that the curvature will be expressed with the aid of the basic covariant differential \bar{D} applied on the covariant components of

the vectors as in (2.2).

Theorem 6. *If in the point M of the curve C in the $R-O_n$ space the mutually orthogonal unit vectors v_0, v_1, \dots, v_{n-1} are constructed so that v_p ($p = 1, \dots, n-1$) satisfies*

$$(2.3) \quad v_{p,i} = \frac{1}{\kappa^{**}} (\tilde{D} v_{p-1,i} + \kappa_{p-1}^{**} v_{p-2,i} + v_{p-1,r} \bar{D} \delta_{i,r}^r)$$

with

$$\kappa_0^{**} = 0, \quad \kappa_n^{**} = 0,$$

then

$$(2.4) \quad \tilde{D} v_{p,i} = P_i^j (-\kappa_{p-1}^{**} v_{p-1,j} + \kappa_{p+1}^{**} v_{p+1,j}) - v_{p,r} D \delta_{i,r}^a Q_a$$

is the Frenet formula with respect to the covariant differential \tilde{D} applied on the covariant components of the observed vectors.

If we make the above calculation with respect to the basic covariant differential \tilde{D} , according to relation (0.9) and the fact that in the case of covariant indices the basic covariant differentials \bar{D} and \tilde{D} are not different, it follows that this case is not different from the observation of Riemannian space. We can only say

Remark 2. *The relation*

$$(2.5) \quad D v_{p,i} = P_i^j (-\kappa_{p-1}^{***} v_{p-1,j} + \kappa_{p+1}^{***} v_{p+1,j})$$

is the Frenet formula with respect to Otsuki's covariant differential D applied on the covariant components of the vectors.

Here $***$, by scalars, denotes that the curvature will be expressed with the aid of Otsuki's basic covariant differential \bar{D} applied on the contravariant components of the vectors.

From (2.5). it follows that

$$(2.6) \quad \kappa_{p+1}^{***} = \left(g^{ij} (\bar{D}v_i + \kappa_{p-1}^{***} v_i) (\bar{D}v_j + \kappa_{p-1}^{***} v_j) \right)^{\frac{1}{2}}$$

holds. Using connection $\bar{D}v_i = \nabla v_i + v_r \bar{D}\delta_i^r$, we get

$$\begin{aligned} \kappa_{p+1}^{***} = & \left(g^{ij} \left(\nabla v_i + \kappa_{p-1}^{***} v_i + v_{pr} \bar{D}\delta_i^r \right) \left(\nabla v_j + \right. \right. \\ & \left. \left. + \kappa_{p-1}^{***} v_j + v_{pq} \bar{D}\delta_j^q \right) \right)^{\frac{1}{2}}. \end{aligned}$$

This means that the value of the curvatures is not different in the case of different differentials. This result can be expected because the curvature depends only on the curve and on a suitably constructed vector frame which is unequally determined. The Frenet formulae are different from the known formulae of Riemannian space only if $D\delta_j^i \neq 0$ holds, i.e. if the basic covariant differential of the metric tensor is not zero.

REFERENCES

- [1] Moór, A.: Otsukische Übertragung mit rekurrentem Masstensor, *Acta Sci. Math. Szeged*, 40 (1978), 129 - 143.
- [2] Moór, A.: Über die durch Kurven bestimmten Sektionalkrümmungen der Finslerschen und Weylschen Räume, *Publ. Math. Debrecen*, 26 (1979), 205 - 214.
- [3] Otsuki, T.: On general connections I, *Math. J. Okayama Univ.*, 9 (1959-60), 99 - 164.
- [4] Tamassy, L.: Frenetschen Formeln für Kurven in affinzusammenhängenden Räumen, *Publ. Math. Debrecen*, 8 (1961).

REZIME

FRENETOVE FORMULE RIMAN-OTSUKIJEVOG PROSTORA

U radu su date Frenetove formule s obzirom na razne kovarijante diferencijale primenjene na kovarijante, odnosno na kontravarijantne indekse vektora.

Received by the editors March 5, 1986.