

CONVERGENCE OF FINITE DIFFERENCE SCHEMES APPROXIMATING HYPERBOLIC PROBLEMS

Boško Jovanović¹, Vladimir Jovanović¹

Abstract. We obtain the convergence rate estimates for the weak solution of hyperbolic initial-boundary value problem of order $2(s - k)/3$, but under conditions weaker than earlier. Also, a numerical example is presented.

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1. Introduction

For the elliptic boundary value problem, convergence rate estimates for finite difference schemes which are compatible with the smoothness of data

$$\|u - v\|_{W_{2,h}^k} \leq Ch^{s-k} \|u\|_{W_2^s}, \quad s > k,$$

has been obtained in [1]. Here u denotes the solution of the boundary value problem, v denotes the corresponding discrete approximation, W_2^s is the standard Sobolev space and $W_{2,h}^k$ is discrete Sobolev space. The compatible estimates may also be derived in parabolic case [2]. But in the hyperbolic case, the usual estimates are not compatible with smoothness of data [3]:

$$\|u - v\|_{C_\tau(W_{2,h}^k)} \leq Ch^{s-k-1} \|u\|_{W_2^s(Q)}, \quad s > k + 1.$$

A few years ago, using interpolation theory (see [4]), B.S. Jovanović derived in [5] the convergence rate estimate of order $2(s - k)/3$ for a finite difference scheme in the case of homogenous wave equation with constant coefficients. The analogous estimate has been obtained in [6] when coefficients are functions from C^k -spaces.

We shall show in this paper that the same estimate holds although the coefficients are less smooth, i.e. if they are from W_2^k -spaces. This will also be illustrated with a numerical experiment.

¹University of Belgrade, Faculty of Mathematics, Studentski trg 16, 11000 Belgrade, Yugoslavia

2. Some necessary function spaces

Let $L_q = L_q(0, 1)$ ($1 \leq q \leq \infty$) be the Lebesgue spaces of integrable functions, $W_2^s = W_2^s(0, 1)$ standard Sobolev spaces, \mathcal{D} the spaces of infinitely differentiable functions with compact support in $(0, 1)$ and $\overset{\circ}{W}_2^s$ be the closure of \mathcal{D} in W_2^s . (\cdot, \cdot) and $\|\cdot\|$ denote the inner product and the norm in L_2 , respectively. Suppose $a \in L_\infty$ such that

$$a \geq a_0 \quad \text{in } (0, 1) \quad \text{a.e.}$$

For the operator $L: \overset{\circ}{W}_2^1 \rightarrow W_2^{-1}$ defined by $Lv = -(av)'$ there exist $0 < \lambda_1 < \lambda_2 < \dots, \lim_k \lambda_k = \infty$ such that $L\varphi_k = \lambda_k \varphi_k$ ($k \in N$); The sequence of the eigenfunctions $(\varphi_k)_{k \in N} \subset \overset{\circ}{W}_2^1$ is an orthonormed basis of L_2 (see [7]). Introduce the spaces V^α ($\alpha \geq 0$) by $V^\alpha = \{v \in L_2 \mid \|v\|_{V^\alpha}^2 = \sum_{k=1}^\infty \lambda_k^\alpha \tilde{v}_k^2 < \infty\}$, where $\tilde{v}_k = (v, \varphi_k)$ are the Fourier coefficients of v in the basis $(\varphi_k)_{k \in N}$.

Lemma 1. *If $v \in L_2$ such that $(\forall \varphi \in \mathcal{D}) \left| \int_0^1 v \varphi' dx \right| \leq K \|\varphi\|$, then $v \in W_2^1$ and $\|v'\| \leq K$.*

Proof. See [8]. □

Lemma 2. *If $v \in W_2^2$ then $\|v'\|_{L_\infty} \leq C \|v\|_{W_2^2}$.*

Proof. We cite here the Poincaré–Wirtinger inequality (see [9])

$$\|v - \bar{v}\|_{L_\infty} \leq \|v'\|_{L_1} \quad (v \in W_1^1),$$

where $\bar{v} = \int_0^1 v dx$ is an average of v . It is easy now to verify our statement by putting v' instead of v in the last equation and applying the Cauchy–Schwarz inequality. □

Lemma 3. *If $v \in \overset{\circ}{W}_2^1$ is the solution of the variational problem*

$$(1) \quad \int_0^1 av' \varphi' dx = \int_0^1 f \varphi dx \quad (\varphi \in \overset{\circ}{W}_2^1),$$

then for $a \in W_2^3$, the relation

$$(2) \quad f \in W_2^i \Rightarrow v \in W_2^{i+2} \quad \text{and} \quad \|v\|_{W_2^{i+2}} \leq C \|f\|_{W_2^i}$$

holds for $i = 0, 1, 2$.

Proof. The case $i = 0, 1$ is simple, but for $i = 2$ the situation is more complicated. So, let us assume that (2) holds for $i = 0, 1$. We shall prove it in the case $i = 2$. Suppose $\varphi \in \mathcal{D}$. Replacing φ with $(\frac{\varphi}{a})''$ in (1) we have

$$\int_0^1 av' \left(\frac{\varphi}{a}\right)''' dx = \int_0^1 f \left(\frac{\varphi}{a}\right)'' dx.$$

Applying partial integration one has

$$\int_0^1 (av')'' \left(\frac{\varphi}{a}\right)' dx = \int_0^1 f'' \frac{\varphi}{a} dx.$$

We know that $v \in W_2^3$ (because (2) holds for $i = 1$), so after a few obvious transformations we obtain

$$\begin{aligned} \int_0^1 v''' \varphi' dx &= \int_0^1 f'' \varphi dx + \int_0^1 \frac{a'}{a^2} (a''v' + 2a'v'' + av''') \varphi dx \\ &- \int_0^1 \frac{a''v' + 2a'v''}{a} \varphi' dx. \end{aligned}$$

Some partial integrations applied to the last equation yield

$$\begin{aligned} \int_0^1 v''' \varphi dx &= \int_0^1 f'' \varphi dx + \int_0^1 \frac{a'}{a} (a''v' + 2a'v'' + av''') \varphi dx + \\ (3) \quad &+ \int_0^1 \left(\frac{a''' + 3a''v'' + 2a'v'''}{a} - \frac{a' a''v' + 2a'v''}{a^2} \right) \varphi dx. \end{aligned}$$

For example, let us estimate the term $\int_0^1 \frac{a' a''}{a^2} v' \varphi dx$. Because $\frac{a' a''}{a^2} \in L_\infty$, we have

$$\begin{aligned} \left| \int_0^1 \frac{a' a''}{a^2} v' \varphi dx \right| &\leq \left\| \frac{a' a''}{a^2} \right\|_{L_\infty} \int_0^1 |v' \varphi| dx \leq C \|v'\| \|\varphi\| \\ &\leq C \|v\|_{W_2^2} \|\varphi\| \leq C \|f\| \|\varphi\| \end{aligned}$$

(the last inequality follows from (2) for $i = 0$).

All the other terms can be estimated similarly except the term $\int_0^1 \frac{a''' v'}{a} \varphi dx$. But $v' \in L_\infty$ and taking into account Lemma 2, we obtain

$$\begin{aligned} \left| \int_0^1 \frac{a''' v'}{a} \varphi dx \right| &\leq \|v'\|_{L_\infty} \int_0^1 \left| \frac{a'''}{a} \varphi \right| dx \leq \|v\|_{W_2^2} \left\| \frac{a'''}{a} \right\| \|\varphi\| \\ &\leq C \|f\| \|\varphi\|. \end{aligned}$$

Estimating all the other terms in the right-hand side of (3) one obtains

$$\left| \int_0^1 v''' \varphi' dx \right| \leq C \|f\|_{W_2^2} \|\varphi\|.$$

Then Lemma 1 implies (2) for $i = 2$. \square

The two following statements can be established in the perfectly same way as it was done in [10].

Theorem 1. *Suppose $a \in W_2^3$ and $a \geq a_0$ in $(0, 1)$. Then*

$$V^1 = \overset{\circ}{W}_2^1, \quad V^2 = W_2^2 \cap \overset{\circ}{W}_2^1, \quad \overset{\circ}{W}_2^3 \subset V^3 \subset W_2^3, \quad \overset{\circ}{W}_2^4 \subset V^4 \subset W_2^4.$$

Proposition 1 *Suppose $a \in W_2^3$ and $a \geq a_0$ in $(0, 1)$. Then*

$$c\|v\|_{W_2^i} \leq \|v\|_{V^i} \leq C\|v\|_{W_2^i} \quad (v \in V^i)$$

for $i = 1, 2, 3, 4$.

3. Hyperbolic problem and discretisation

Consider the initial-boundary value problem for the homogenous second-order hyperbolic equation in the domain $Q = (0, 1) \times (0, T]$:

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial x} \left(a(x) \frac{\partial u}{\partial x} \right), \quad (x, t) \in Q$$

(4)
$$u(0, t) = u(1, t) = 0, \quad t \in [0, T]$$

$$u(x, 0) = u_0(x), \quad \frac{\partial u}{\partial t}(x, 0) = u_1(x), \quad x \in (0, 1).$$

There is the unique weak solution of this problem for $u_0 \in V^1$, $u_1 \in V^0$ (see [11]). In [6] was also shown that if $u_0 \in V^\alpha$, $u_1 \in V^{\alpha-1}$ and $a \in L_\infty$, $a \geq a_0 > 0$ in $(0, 1)$ then

$$\max_{t \in [-T, T]} \|\partial^l u / \partial t^l\|_{V^{\alpha-l}} \leq C(\|u_0\|_{V^\alpha} + \|u_1\|_{V^{\alpha-1}}),$$

for $1 \leq \alpha \leq 4$, $l \in Z$, $0 \leq l \leq \alpha$, wherefrom, using Theorem 1 and Proposition 1 follows that if $a \in W_2^3$ and $a \geq a_0 > 0$ in $(0, 1)$ then

(5)
$$\max_{t \in [-T, T]} \|\partial^l u / \partial t^l\|_{W_2^{\alpha-l}} \leq C(\|u_0\|_{V^\alpha} + \|u_1\|_{V^{\alpha-1}}),$$

for $1 \leq \alpha \leq 4$, $l \in Z$, $0 \leq l \leq \alpha$. The extension of $\partial^l u / \partial t^l$ on $[-T, T]$ can be done in a natural way using the Fourier series (see also [6]).

Let $\bar{\omega}_h$ be a uniform mesh on $[0, 1]$ with the stepsize $h = 1/n$, $\omega_h = \bar{\omega}_h \cap (0, 1)$ and $\omega^- = \omega_h \cup \{0\}$. We set $\overset{\circ}{W}_{2,h}^1$ to be the space of all functions defined on $\bar{\omega}_h$ vanishing at 0 and 1. Introduce finite differences in x :

$$v_x = (v(x+h) - v(x))/h, \quad v_{\bar{x}} = (v(x) - v(x-h))/h.$$

We define the following discrete norms

$$\|v\|_h = \left(h \sum_{x \in \omega_h} v^2(x) \right)^{1/2}, \quad \|v\|_h = \left(h \sum_{x \in \omega_h^-} v^2(x) \right)^{1/2},$$

$$\|v\|_{W_{2,h}^1} = (\|v\|^2 + \|v_x\|^2)^{1/2}.$$

Also introduce the operator $L_h : \overset{\circ}{W}_{2,h}^1 \rightarrow \overset{\circ}{W}_{2,h}^1$ by

$$L_h v = \begin{cases} -\frac{1}{2}[(av_x)_{\bar{x}} + (av_{\bar{x}})_x], & x \in \omega_h \\ 0, & x \in \{0, 1\} \end{cases}$$

Let $\bar{\omega}_h$ be a uniform mesh on $[-\tau/2, T]$ with the stepsize $\tau = T/(m-1/2)$, $\omega_\tau = \bar{\omega}_\tau \cap (0, T)$, and $\omega_\tau^- = \omega \cup \{-\tau/2\}$. We shall also use the following notations:

$$v = v(t), \quad \hat{v} = v(t + \tau), \quad \check{v} = v(t - \tau), \quad v^j = v((j - 1/2)\tau),$$

$$\bar{v} = (v + \hat{v})/2, \quad v_t = (\hat{v} - v)/\tau, \quad v_{\bar{t}} = (v - \check{v})/\tau.$$

For the functions defined on $\bar{\omega}_h \times \bar{\omega}_\tau$ we define the norms

$$\|v\|_{C_\tau(W_{2,h}^1)} = \max_{t \in \omega_\tau^-} \|v(\cdot, t)\|_{W_{2,h}^1}, \quad \|v\|_{L_{q,\tau}(L_{2,h})} = \left(\tau \sum_{t \in \omega_\tau} \|v(\cdot, t)\|_h^q \right)^{1/q}.$$

One can easily deduce

Lemma 4. For $v \in \overset{\circ}{W}_{2,h}^1$ the inequality $\|v\|_{(I+0.25\tau^2(\sigma-1/4)L_h)} \leq C\|v\|_h$ holds if one of the following two conditions is satisfied:

- (i) If $\sigma > 1/4$, then $\tau/h < C$, where C is an arbitrary constant;
- (ii) If $\sigma < 1/4$, then $\tau/h \leq 4\sqrt{\frac{1-s_0}{(1-4\sigma)c_1^2}}$ for some $s_0 \in (0, 1)$, where c_1 is a constant depending only on the function a .

Let S_x and S_t denote the Steklov smoothing operators in x and t :

$$S_x f(x, t) = \frac{1}{h} \int_{x-h/2}^{x+h/2} f(s, t) ds, \quad S_t f(x, t) = \frac{1}{\tau} \int_{t-\tau/2}^{t+\tau/2} f(x, \eta) d\eta.$$

For the approximation of (4) we shall use a weighted finite difference scheme:

$$v_{\bar{t}\bar{t}} = -L_h(\sigma\hat{v} + (1 - 2\sigma)v + \sigma\check{v}),$$

$$v(0, t) = v(1, t) = 0, \quad t \in \bar{\omega}_\tau,$$

$$v^0 = u_0 - (\tau/2)S_x^2 u_1, \quad v^1 = u_0 + (\tau/2)S_x^2 u_1.$$

Finally, we have

Theorem 2. Suppose $a \in W_2^3$ and $a \geq a_0$ in $(0, 1)$, u is the weak solution of (4), v is its discrete approximation and let one of the conditions in Lemma 4 is satisfied. Then for the error $z = u - v$ the following estimates hold:

$$(i) \quad \|\bar{z}\|_{C_\tau(W_{2,h}^1)} \leq C(h + \tau)^{2(s-1)/3} (\|u_0\|_{V^s} + \|u_1\|_{V^{s-1}}),$$

$$\text{if } u_0 \in V^s, u_1 \in V^{s-1}, 1 \leq s \leq 4;$$

$$(ii) \quad \|\bar{z}\|_{C_\tau(W_{2,h}^1)} \leq C(h + \tau)^{2(s-1)/3} (\|u_0\|_{W_2^s} + \|u_1\|_{W_2^{s-1}}),$$

$$\text{if } u_0 \in \overset{\circ}{W}_2^s, u_1 \in \overset{\circ}{W}_2^{s-1}, 1 \leq s \leq 4, s \neq \text{integer} + 1/2.$$

Proof. Starting with (5) and applying identical techniques that we used for deriving the analogous estimates in [6], we prove this theorem. \square

4. Numerical examples

A series of initial-boundary value problems of the form

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial x} \left(a(x) \frac{\partial u}{\partial x} \right), \quad (x, t) \in Q = (0, 1) \times (0, T).$$

$$u(0, t) = u(1, t) = 0, \quad u(x, 0) = u_0(x), \quad \frac{\partial u(x, 0)}{\partial x} = 0$$

were solved numerically. Input data were taken in the form

$$u_0(x) = \psi_{s_0}(x), \quad a(x) = 1 + \psi_{s_1}(x),$$

where

$$\psi_s(x) = \begin{cases} 0, & x \in (0, 1/2) \\ 4^{s+2.5} (x - 0.5)^{s-0.5} (1-x)^3, & x \in (1/2, 1). \end{cases}$$

We have $\psi_s(x) \in W_2^{s-\varepsilon}(0, 1)$ for all $\varepsilon > 0$.

The problem was solved for various values of parameters s_0 and s_1 , using symmetric difference scheme with the weight $\sigma = 1/4$, on the sequence of refined meshes. The rate of convergence was determined numerically, using Runge's rule.

Obtained convergence rates for $T = 1/2$, $s_1 = 3$ and $1 \leq s_0 \leq 4$ are presented in the following table.

Mesh	s_0						
	4.0	3.5	3.0	2.5	2.0	1.5	1.0
16 × 16	-0.163	-0.221	-0.226	-0.211	-0.335	-0.660	-1.375
32 × 32	0.510	0.714	0.798	0.877	0.934	0.341	-0.165
64 × 64	1.481	1.553	1.674	1.542	0.860	0.492	0.028
128 × 128	1.861	1.831	1.640	1.178	0.755	0.244	-0.129
256 × 256	1.951	1.890	1.589	0.958	0.423	0.148	-0.188
512 × 512	1.994	1.889	1.434	1.004	0.662	0.307	-0.097
$\frac{2}{3}(s_0 - 1)$	2.000	1.667	1.333	1.000	0.667	0.333	0.000

We note that the numerical convergence rates weakly depend on the smoothness of the coefficient $a(x)$. The obtained convergence rates for $T = 1/2$, $s_0 = 4$ and $0.5 \leq s_1 \leq 3$ are presented in the following table.

Mesh	s_1					
	3.0	2.5	2.0	1.5	1.0	0.5
16 × 16	-0.163	-0.254	-0.286	-0.142	-0.146	-0.094
32 × 32	0.510	0.619	0.653	0.503	0.351	0.134
64 × 64	1.481	1.466	1.375	1.141	0.917	1.103
128 × 128	1.861	1.869	1.831	1.706	1.410	1.061
256 × 256	1.951	1.960	1.965	1.915	1.679	1.138
512 × 512	1.994	2.006	2.009	2.002	1.837	1.398

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