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SOME CHARACTERIZATIONS OF INCLINED CURVES IN EUCLIDEAN E^n SPACE

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Abstract. We consider a unit speed curve α in the Euclidean *n*-dimensional space \mathbf{E}^n and denote the Frenet frame of α by $\{\mathbf{V}_1, \ldots, \mathbf{V}_n\}$. We say that α is a cylindrical helix if its tangent vector \mathbf{V}_1 makes a constant angle with a fixed direction U. In this work we give different characterizations of such curves in terms of their curvatures.

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1. Introduction and statement of results

An helix in the Euclidean 3-space \mathbf{E}^3 is a curve where the tangent lines make a constant angle with a fixed direction. An helix curve is characterized by the fact that the ratio κ/τ is constant along the curve, where κ and τ are the curvature and the torsion of α , respectively. Helices are well known curves in classical differential geometry of space curves [4] and we refer to the reader for recent works on this type of curves [2, 7]. Recently, Magden [3] have introduced the concept of cylindrical helix in the Euclidean 4-space \mathbf{E}^4 , saying that the tangent lines make a constant angle with a fixed directions. He characterizes a cylindrical helix in \mathbf{E}^4 if and only if the function

(1)
$$\left(\frac{\kappa_1}{\kappa_2}\right)^2 + \left(\frac{1}{\kappa_3}\left(\frac{\kappa_1}{\kappa_2}\right)'\right)^2$$

is constant along the curve, where κ_3 and κ_4 are the third and the fourth curvature of the the curve. See also [5].

In this work we consider the generalization of the concept of general helices in the Euclidean n-space \mathbf{E}^n . Let $\alpha : I \subset \mathbb{R} \to \mathbf{E}^n$ be an arbitrary curve in \mathbf{E}^n . Recall that the curve α is said to be of unit speed (or parameterized by

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the arc-length function s) if $\langle \alpha'(s), \alpha'(s) \rangle = 1$, where \langle , \rangle is the standard scalar product in the Euclidean space \mathbf{E}^n given by

$$\langle X, Y \rangle = \sum_{i=1}^{n} x_i y_i,$$

for each $X = (x_1, ..., x_n), Y = (y_1, ..., y_n) \in \mathbf{E}^n$.

Let $\{\mathbf{V}_1(s), \ldots, \mathbf{V}_n(s)\}$ be the moving frame along α , where the vectors \mathbf{V}_i are mutually orthogonal vectors satisfying $\langle \mathbf{V}_i, \mathbf{V}_i \rangle = 1$. The Frenet equations for α are given by ([2])

$$(2) \begin{bmatrix} \mathbf{V}_{1}' \\ \mathbf{V}_{2}' \\ \mathbf{V}_{3}' \\ \vdots \\ \mathbf{V}_{n-1}' \\ \mathbf{V}_{n}' \end{bmatrix} = \begin{bmatrix} 0 & \kappa_{1} & 0 & 0 & \cdots & 0 & 0 \\ -\kappa_{1} & 0 & \kappa_{2} & 0 & \cdots & 0 & 0 \\ 0 & -\kappa_{2} & 0 & \kappa_{3} & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \cdots & 0 & \kappa_{n-1} \\ 0 & 0 & 0 & 0 & \cdots & -\kappa_{n-1} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{V}_{1} \\ \mathbf{V}_{2} \\ \mathbf{V}_{3} \\ \vdots \\ \mathbf{V}_{n-1} \\ \mathbf{V}_{n} \end{bmatrix}.$$

Recall that the functions $\kappa_i(s)$ are called the i-th curvatures of α . If $\kappa_{n-1}(s) = 0$ for any $s \in I$, then $\mathbf{V}_n(s)$ is a constant vector V and the curve α lies in an (n-1)-dimensional affine subspace orthogonal to V, which is isometric to the Euclidean (n-1)-space \mathbf{E}^{n-1} . We will assume throughout this work that all the curvatures satisfy $\kappa_i(s) \neq 0$ for any $s \in I$, $1 \leq i \leq n-1$.

Definition 1.1. A unit speed curve $\alpha : I \to \mathbf{E}^n$ is called cylindrical helix if its tangent vector \mathbf{V}_1 makes a constant angle with a fixed direction U.

Our main result in this work is the following characterization of cylindrical helices in the Euclidean *n*-space \mathbf{E}^{n} .

Theorem 1.2. Let $\alpha : I \to \mathbf{E}^n$ be a unit speed curve in \mathbf{E}^n . Define the functions

(3)
$$G_1 = 1, \ G_2 = 0, \ G_i = \frac{1}{\kappa_{i-1}} \Big[\kappa_{i-2} G_{i-2} + G'_{i-1} \Big], \ 3 \le i \le n.$$

Then α is a cylindrical helix if and only if the function

(4)
$$\sum_{i=3}^{n} G_i^2 = C$$

is constant. Moreover, the constant $C = \tan^2 \theta$, θ being the angle that makes \mathbf{V}_1 with the fixed direction U that determines α .

This theorem generalizes in arbitrary dimensions what happens for n = 3and n = 4, namely: if n = 3, (4) writes $G_3^2 = \kappa_1/\kappa_2 = \kappa/\tau$ and for n = 4, (4) agrees with (1).

2. Proof of Theorem 1.2

Let α be a unit speed curve in \mathbf{E}^n . Assume that α is a cylindrical helix curve. Let U be the direction with which \mathbf{V}_1 makes a constant angle θ and, without loss of generality, we suppose that $\langle U, U \rangle = 1$. Consider the differentiable functions $a_i, 1 \leq i \leq n$,

(5)
$$U = \sum_{i=1}^{n} a_i(s) \mathbf{V}_i(s), \quad s \in I,$$

that is,

$$a_i = \langle \mathbf{V}_i, U \rangle, \ 1 \le i \le n.$$

Then the function $a_1(s) = \langle \mathbf{V}_1(s), U \rangle$ is constant, and it agrees with $\cos \theta$:

(6)
$$a_1(s) = \langle \mathbf{V}_1, U \rangle = \cos \theta$$

for any s. By differentiating (6) with respect to s and using the Frenet formula (2) we have

$$a_1'(s) = \kappa_1 \left\langle \mathbf{V}_2, U \right\rangle = \kappa_1 \, a_2 = 0.$$

Then $a_2 = 0$ and therefore U lies in the subspace $Sp(\mathbf{V}_1, \mathbf{V}_3, \dots, \mathbf{V}_n)$. Because the vector field U is constant, a differentiation in (5), together with (2) gives the following ordinary differential equation system

Define the functions $G_i = G_i(s)$ as follows

(8)
$$a_i(s) = G_i(s) a_1, \ 3 \le i \le n.$$

We point out that $a_1 \neq 0$: on the contrary, (7) gives $a_i = 0$, for $3 \leq i \leq n$ and so, U = 0: contradiction. The first (n - 2)-equations in (7) lead to

$$G_{3} = \frac{\kappa_{1}}{\kappa_{2}}$$

$$G_{4} = \frac{1}{\kappa_{3}}G'_{3}$$

$$G_{5} = \frac{1}{\kappa_{4}} \left[\kappa_{3}G_{3} + G'_{4} \right]$$

$$\vdots$$

$$G_{n-1} = \frac{1}{\kappa_{n-2}} \left[\kappa_{n-3}G_{n-3} + G'_{n-2} \right]$$

$$G_{n} = \frac{1}{\kappa_{n-1}} \left[\kappa_{n-2}G_{n-2} + G'_{n-1} \right]$$

(9)

The last equation of (7) leads to the following condition;

(10)
$$G'_n + \kappa_{n-1} G_{n-1} = 0.$$

We do the change of variables:

$$t(s) = \int^{s} \kappa_{n-1}(u) du, \quad \frac{dt}{ds} = \kappa_{n-1}(s).$$

In particular, from the last equation of (9), we have

$$G'_{n-1}(t) = G_n(t) - \left(\frac{\kappa_{n-2}(t)}{\kappa_{n-1}(t)}\right)G_{n-2}(t).$$

As a consequence, if α is a cylindrical helix, substituting the equation (10) in the last equation yields

$$G_n''(t) + G_n(t) = \frac{\kappa_{n-2}(t)G_{n-2}(t)}{\kappa_{n-1}(t)}.$$

The general solution of this equation is

(11)
$$G_{n}(t) = \left(A - \int \frac{\kappa_{n-2}(t)G_{n-2}(t)}{\kappa_{n-1}(t)}\sin t \, dt\right)\cos t + \left(B + \int \frac{\kappa_{n-2}(t)G_{n-2}(t)}{\kappa_{n-1}(t)}\cos t \, dt\right)\sin t,$$

where A and B are arbitrary constants. Then (11) takes the following form (12)

$$G_n(s) = \left(A - \int \left[\kappa_{n-2}(s)G_{n-2}(s)\sin\int\kappa_{n-1}(s)ds\right]ds\right)\cos\int\kappa_{n-1}(s)ds + \left(B + \int \left[\kappa_{n-2}(s)G_{n-2}(s)\cos\int\kappa_{n-1}(s)ds\right]ds\right)\sin\int\kappa_{n-1}(s)ds.$$

From (10), the function G_{n-1} is given by (13)

$$G_{n-1}(s) = \left(A - \int \left[\kappa_{n-2}(s)G_{n-2}(s)\sin\int\kappa_{n-1}(s)ds\right]ds\right)\sin\int\kappa_{n-1}(s)ds - \left(B + \int \left[\kappa_{n-2}(s)G_{n-2}(s)\cos\int\kappa_{n-1}(s)ds\right]ds\right)\cos\int\kappa_{n-1}(s)ds.$$

From equation (9), we have

$$\sum_{i=3}^{n-2} G_i G'_i = G_3 \kappa_3 G_4 + G_4 \left(\kappa_4 G_5 - \kappa_3 G_3 \right) + \dots$$

+ $G_{n-3} \left(\kappa_{n-3} G_{n-2} - \kappa_{n-4} G_{n-4} \right) + G_{n-2} G'_{n-2}$
= $G_{n-2} \left(G'_{n-2} + \kappa_{n-3} G_{n-3} \right)$
= $\kappa_{n-2} G_{n-2} G_{n-1}$

Substituting (13) in the above equation and integrate it, we have

(14)
$$\sum_{i=3}^{n-2} G_i^2 = C - \left(A - \int \left[\kappa_{n-2}(s)G_{n-2}(s)\sin\int\kappa_{n-1}ds\right]ds\right)^2 - \left(B + \int \left[\kappa_{n-2}(s)G_{n-2}(s)\cos\int\kappa_{n-1}ds\right]ds\right)^2,$$

where C is a constant of integration. From Equations (12) and (13), we have

(15)
$$G_n^2 + G_{n-1}^2 = \left(A - \int \left[\kappa_{n-2}(s)G_{n-2}(s)\sin\int\kappa_{n-1}ds\right]ds\right)^2 + \left(B + \int \left[\kappa_{n-2}(s)G_{n-2}(s)\cos\int\kappa_{n-1}ds\right]ds\right)^2,$$

It follows from (14) and (15) that

$$\sum_{i=3}^{n} G_i^2 = C.$$

Moreover, the constant C is calculated as follows. From (8), together with the (n-2)-equations (9), we have

$$C = \sum_{i=3}^{n} G_i^2 = \frac{1}{a_1^2} \sum_{i=3}^{n} a_i^2 = \frac{1 - a_1^2}{a_1^2} = \tan^2 \theta,$$

where we have used (2) and the fact that U is a unit vector field.

We do the converse of Theorem. Assume that the condition (9) is satisfied for a curve α . Let $\theta \in \mathbb{R}$ be so that $C = \tan^2 \theta$. Define the unit vector U by

$$U = \cos\theta \Big[\mathbf{V}_1 + \sum_{i=3}^n G_i \, \mathbf{V}_i \Big].$$

By taking into account (9), a differentiation of U gives that $\frac{dU}{ds} = 0$, which means that U is a constant vector field. On the other hand, the scalar product between the unit tangent vector field \mathbf{V}_1 with U is

$$\langle \mathbf{V}_1(s), U \rangle = \cos \theta.$$

Thus α is a cylindrical helix curve. This finishes the proof of Theorem 1.2.

As a direct consequence of the proof, we generalize Theorem 1.2 in Minkowski space and for timelike curves.

Theorem 2.1. Let \mathbf{E}_1^n be the Minkowski n-dimensional space and let $\alpha : I \to \mathbf{E}_1^n$ be a unit speed timelike curve. Then α is a cylindrical helix if and only if the function $\sum_{i=3}^n G_i^2$ is constant, where the functions G_i are defined as in (3).

Proof. The proof is carried the same steps as above, and we omit the details. We only point out that the fact that α is timelike means that $\mathbf{V}_1(s) = \alpha'(s)$ is a timelike vector field. The other V_i in the Frenet frame, $2 \leq i \leq n$, are unit spacelike vectors, and the second equation in (2) changes to $\mathbf{V}'_2 = \kappa_1 \mathbf{V}_1 + \kappa_2 \mathbf{V}_3$ ([1, 6]).

3. Further characterizations of cylindrical helices

In this section we present new characterizations of a cylindrical helix in \mathbf{E}^{n} . The first one is a consequence of Theorem 1.2.

Theorem 3.1. Let $\alpha : I \subset R \to \mathbf{E}^n$ be a unit speed curve in the Euclidean space \mathbf{E}^n . Then α is a cylindrical helix if and only if there exists a C^2 -function $G_n(s)$ such that

(16)
$$G_n = \frac{1}{\kappa_{n-1}} \Big[\kappa_{n-2} G_{n-2} + G'_{n-1} \Big], \quad \frac{dG_n}{ds} = -\kappa_{n-1}(s) G_{n-1}(s),$$

where

$$G_1 = 1, G_2 = 0, G_i = \frac{1}{\kappa_{i-1}} \left[\kappa_{i-2} G_{i-2} + G'_{i-1} \right], \ 3 \le i \le n-1.$$

Proof. Let now assume that α is a cylindrical helix. By using Theorem 1.2 and by the differentiation of the (constant) function given in (4), we obtain

$$0 = \sum_{i=3}^{n} G_{i} G'_{i}$$

= $G_{3}\kappa_{3}G_{4} + G_{4} (\kappa_{4}G_{5} - \kappa_{3}G_{3}) + \dots$
 $\dots + G_{n-1} (\kappa_{n-1}G_{n} - \kappa_{n-2}G_{n-2}) + G_{n}G'_{n}$
= $G_{n} (G'_{n} + \kappa_{n-1}G_{n-1}).$

This shows (16). Conversely, if (16) holds, we define a vector field U by

$$U = \cos\theta \Big[\mathbf{V}_1 + \sum_{i=3}^n G_i \, \mathbf{V}_i \Big].$$

By the Frenet equations (2), $\frac{dU}{ds} = 0$, and so, U is constant. On the other hand, $\langle \mathbf{V}_1(s), U \rangle = \cos \theta$ is constant, and this means that α is a cylindrical helix. \Box

At the end, we give an integral characterization of a cylindrical helix.

Theorem 3.2. Let $\alpha : I \subset R \to \mathbf{E}^n$ be a unit speed curve in the Euclidean space \mathbf{E}^n . Then α is a cylindrical helix if and only if the following condition is satisfied

(17)
$$G_{n-1}(s) = \left(A - \int \left[\kappa_{n-2}G_{n-2}\sin\int\kappa_{n-1}ds\right]ds\right)\sin\int^{s}\kappa_{n-1}(u)du \\ -\left(B + \int \left[\kappa_{n-2}G_{n-2}\cos\int\kappa_{n-1}ds\right]ds\right)\cos\int^{s}\kappa_{n-1}(u)du.$$

for some constants A and B.

Proof. Suppose that α is a cylindrical helix. By using Theorem 3.1, let define m(s) and n(s) by

$$\phi(s) = \int^s \kappa_{n-1}(u) du$$

(18)
$$m(s) = G_n(s)\cos\phi + G_{n-1}(s)\sin\phi + \int \kappa_{n-2}G_{n-2}\sin\phi \,ds, n(s) = G_n(s)\sin\phi - G_{n-1}(s)\cos\phi - \int \kappa_{n-2}G_{n-2}\cos\phi \,ds.$$

If we differentiate equations (18) with respect to s, and taking into account of (17) and (16), we obtain $\frac{dm}{ds} = 0$ and $\frac{dn}{ds} = 0$. Therefore, there exist constants A and B such that m(s) = A and n(s) = B. By substituting into (18) and solving the resulting equations for $G_{n-1}(s)$, we get

$$G_{n-1}(s) = \left(A - \int \kappa_{n-2} G_{n-2} \sin \phi \, ds\right) \sin \phi - \left(B + \int \kappa_{n-2} G_{n-2} \cos \phi \, ds\right) \cos \phi.$$

Conversely, suppose that (17) holds. In order to apply Theorem 3.1, we define $G_n(s)$ by

$$G_n(s) = \left(A - \int \kappa_{n-2} G_{n-2} \sin \phi \, ds\right) \cos \phi + \left(B + \int \kappa_{n-2} G_{n-2} \cos \phi \, ds\right) \sin \phi.$$

with $\phi(s) = \int^{s} \kappa_{n-1}(u) du$. A direct differentiation of (17) gives

$$G_{n-1}' = \kappa_{n-1}G_n - \kappa_{n-2}G_{n-2}.$$

This shows the left condition in (16). Moreover, a straightforward computation leads to $G'_n(s) = -\kappa_{n-1}G_{n-1}$, which finishes the proof.

We end this section with a characterization of cylindrical helices only in terms of the curvatures of α . From the definitions of G_i in (3), one can express the functions G_i in terms of G_3 and the curvatures of α as follows:

(19)
$$G_j = \sum_{i=0}^{j-3} A_{ji} G_3^{(i)}, \ 3 \le j \le n,$$

where

$$G_3^{(i)} = \frac{d^{(i)}G_3}{ds^i}, \ G_3^{(0)} = G_3 = \frac{\kappa_1}{\kappa_2},$$

Then

$$G_4 = \kappa_3^{-1}G'_3 = A_{41}G'_3 + A_{40}G_3, A_{41} = \kappa_3^{-1}, A_{40} = 0$$

$$G_5 = A_{52}G''_3 + A_{51}G'_3 + A_{50}G_3, A_{52} = \kappa_4^{-1}A_{41}, A_{51} = \kappa_4^{-1}A'_{41}, A_{50} = \kappa_4^{-1}\kappa_3$$

and so on. Define the functions $A_{ji} = A_{ij}(s)$, $3 \le j$, $0 \le i \le j-3$ as the following:

$$A_{30} = 1, A_{40} = 0$$
$$A_{j0} = \kappa_{j-1}^{-1} \kappa_{j-2} A_{(j-2)0} + \kappa_{j-1}^{-1} A'_{(j-1)0}, \ 5 \le j \le n$$

$$\begin{aligned} A_{j(j-3)} &= \kappa_{j-1}^{-1} \kappa_{j-2}^{-1} \kappa_{j-3}^{-1} \dots \kappa_{4}^{-1} \kappa_{3}^{-1}, \text{for } 4 \leq j \leq n \\ A_{j(j-4)} &= \kappa_{j-1}^{-1} \left(\kappa_{j-2}^{-1} \kappa_{j-3}^{-1} \dots \kappa_{4}^{-1} \kappa_{3}^{-1}\right)' + \kappa_{j-1}^{-1} \kappa_{j-2}^{-1} \left(\kappa_{j-3}^{-1} \dots \kappa_{4}^{-1} \kappa_{3}^{-1}\right)' \\ &+ \dots + \kappa_{j-1}^{-1} \kappa_{j-2}^{-1} \kappa_{j-3}^{-1} \dots \kappa_{4}^{-1} \left(\kappa_{3}^{-1}\right)', \\ &\text{for } 5 \leq j \leq n \\ A_{ji} &= \kappa_{j-1}^{-1} \kappa_{j-2} A_{(j-2)i} + \kappa_{j-1}^{-1} \left(A'_{(j-1)i} + A_{(j-1)(i-1)}\right) \\ &\text{for } 1 \leq i \leq j - 5, 6 \leq j \leq n \end{aligned}$$

and $A_{ji} = 0$ otherwise.

The second equation of (16) leads the following condition:

(20)

$$\begin{array}{rcl}
A_{n(n-3)}G_{3}^{(n-2)} &+ \left(A_{n(n-3)}' + A_{n(n-4)}\right)G_{3}^{(n-3)} \\
&+ \sum_{i=1}^{n-4} \left[A_{ni}' + A_{n(i-1)} + \kappa_{n-1}A_{(n-1)i}\right]G_{3}^{(i)} \\
&+ \left(A_{n0}' + \kappa_{n-1}A_{(n-1)0}\right)G_{3} = 0, \quad n \ge 3.
\end{array}$$

As a consequence of (20) and Theorem 1.2, we have the following corollary.

Corollary 3.3. Let $\alpha : I \to \mathbf{E}^n$ be a unit speed curve in \mathbf{E}^n . The next statements are equivalent:

- 1. α is a cylindrical helix.
- 2.

$$0 = A_{n(n-3)} \left(\frac{\kappa_1}{\kappa_2}\right)^{(n-2)} + \left(A'_{n(n-3)} + A_{n(n-4)}\right) \left(\frac{\kappa_1}{\kappa_2}\right)^{(n-3)} + \sum_{i=1}^{n-4} \left[A'_{ni} + A_{n(i-1)} + \kappa_{n-1}A_{(n-1)i}\right] \left(\frac{\kappa_1}{\kappa_2}\right)^{(i)} + \left(A'_{n0} + \kappa_{n-1}A_{(n-1)0}\right) \left(\frac{\kappa_1}{\kappa_2}\right), \quad n \ge 3.$$

3. The function

$$\sum_{j=3}^{n} \sum_{i=0}^{j-3} \sum_{k=0}^{j-3} A_{ji} A_{jk} \left(\frac{\kappa_1}{\kappa_2}\right)^{(i)} \left(\frac{\kappa_1}{\kappa_2}\right)^{(k)} = C$$

is constant, $j - i \ge 3$, $j - k \ge 3$.

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