# APPLICATIONS OF THE THICK DISTRIBUTIONAL CALCULUS 

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#### Abstract

We give several applications of the thick distributional calculus. We consider homogeneous thick distributions, point source fields, and higher order derivatives of order 0 .


AMS Mathematics Subject Classification (2010): 46F10
Key words and phrases: thick points, delta functions, distributions, generalized functions

## 1. Introduction

The aim of this note is to give several applications of the recently introduced calculus of thick distributions in several variables [16, generalizing the thick distributions of one variable [3]. The thick distributional calculus allows us to study problems where a finite number of special points are present; it is the distributional version of the analysis of Blanchet and Faye [1], who employed the concepts of Hadamard finite parts as developed by Sellier [13] to study dynamics of point particles in high post-Newtonian approximations of general relativity. We give a short summary of the theory of thick distributions in Section 2

Our first application, given in Section 3, is the computation of the distributional derivatives of homogeneous distributions in $\mathbb{R}^{n}$ by first computing the thick distributional derivatives and then projecting onto the space of standard distributions. Our analysis makes several delicate points quite clear.

Next, in Section 4, we consider an application to point source fields. In [2], Bowen computed the derivative of the distribution

$$
\begin{equation*}
g_{j_{1}, \ldots, j_{k}}(\mathbf{x})=\frac{n_{j_{1}} \cdots n_{j_{k}}}{r^{2}} \tag{1.1}
\end{equation*}
$$

of $\mathcal{D}^{\prime}\left(\mathbb{R}^{3}\right)$, where $r=|\mathbf{x}|$ and $\mathbf{n}=\left(n_{i}\right)$ is the unit normal vector to a sphere centered at the origin, that is, $n_{i}=x_{i} / r$. Following the notation introduced by the late Professor Farassat [8] of denoting distributional derivatives with an overbar, Bowen's result can be written as,
(1.2) $\frac{\bar{\partial}}{\partial x_{i}} g_{j_{1}, \ldots, j_{k}}=\left\{\sum_{q=1}^{k} \delta_{i j_{q}} \frac{n_{j_{1}} \cdots n_{j_{k}}}{n_{j_{q}}}-(k+2) n_{i} n_{j_{1}} \cdots n_{j_{k}}\right\} \frac{1}{r^{3}}+A \delta(\mathbf{x})$,

[^0]where $n_{i} n_{j_{1}} \cdots n_{j_{k}}=n_{1}^{a} n_{2}^{b} n_{3}^{c}$, and $A=0$ if $a, b$, or $c$ is odd, while
\[

$$
\begin{equation*}
A=\frac{2 \Gamma((a+1) / 2) \Gamma((b+1) / 2) \Gamma((c+1) / 2)}{\Gamma((a+b+c+3) / 2)} \tag{1.3}
\end{equation*}
$$

\]

if the three exponents are even. Interestingly, he observes that if one tries to compute this formula by induction, employing the product rule for derivatives, the result obtained is wrong. In this article we show that one can actually apply the product rule in the space of thick distributions, obtaining $\sqrt{1.2}$ by induction; furthermore, our analysis shows why the wrong result is obtained when applying the product rule in [2].

Finally in Section 5 we show how the thick distributional calculus allows one to avoid mistakes in the computation of higher order derivatives of thick distributions of order 0 .

## 2. Thick distributions

We now recall the basic ideas of the thick distributional calculus [16]. If a is a fixed point of $\mathbb{R}^{n}$, then the space of test functions with a thick point at $\mathbf{x}=\mathbf{a}$ is defined as follows.

Definition 2.1. Let $\mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right)$ denote the vector space of all smooth functions $\phi$ defined in $\mathbb{R}^{n} \backslash\{\mathbf{a}\}$, with support of the form $K \backslash\{\mathbf{a}\}$, where $K$ is compact in $\mathbb{R}^{n}$, that admit a strong asymptotic expansion of the form

$$
\begin{equation*}
\phi(\mathbf{a}+\mathbf{x})=\phi(\mathbf{a}+r \mathbf{w}) \sim \sum_{j=m}^{\infty} a_{j}(\mathbf{w}) r^{j}, \quad \text { as } \mathbf{x} \rightarrow \mathbf{0} \tag{2.1}
\end{equation*}
$$

where $m$ is an integer (positive or negative), and where the $a_{j}$ are smooth functions of $\mathbf{w}$, that is, $a_{j} \in \mathcal{D}(\mathbb{S})$. The subspace $\mathcal{D}_{*, \mathbf{a}}^{[m]}\left(\mathbb{R}^{n}\right)$ consists of those test functions $\phi$ whose expansion (2.1) begins at $m$. For a fixed compact $K$ whose interior contains a, $\mathcal{D}_{*, \mathbf{a}}^{[m ; K]}\left(\mathbb{R}^{n}\right)$ is the subspace formed by those test functions of $\mathcal{D}_{*, \mathbf{a}}^{[m]}\left(\mathbb{R}^{n}\right)$ that vanish in $\mathbb{R}^{n} \backslash K$.

Observe that we require the asymptotic development of $\phi(\mathbf{x})$ as $\mathbf{x} \rightarrow$ a to be "strong". This means [7, Chapter 1] that for any differentiation operator $(\partial / \partial \mathbf{x})^{\mathbf{p}}=\left(\partial^{p_{1}} \ldots \partial^{p_{n}}\right) / \partial x_{1}^{p_{1}} \ldots \partial x_{n}^{p_{n}}$, the asymptotic development of $(\partial / \partial \mathbf{x})^{\mathbf{p}} \phi(\mathbf{x})$ as $\mathbf{x} \rightarrow \mathbf{a}$ exists and is equal to the term-by-term differentiation of $\sum_{j=m}^{\infty} a_{j}(\mathbf{w}) r^{j}$. Observe that saying that the expansion exists as $\mathbf{x} \rightarrow \mathbf{0}$ is the same as saying that it exists as $r \rightarrow 0$, uniformly with respect to $\mathbf{w}$.

We call $\mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right)$ the space of test functions on $\mathbb{R}^{n}$ with a thick point located at $\mathbf{x}=\mathbf{a}$. We denote $\mathcal{D}_{*, \mathbf{0}}\left(\mathbb{R}^{n}\right)$ as $\mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$.

The topology of the space of thick test functions is constructed as follows.
Definition 2.2. Let $m$ be a fixed integer and $K$ a compact subset of $\mathbb{R}^{n}$ whose interior contains a. The topology of $\mathcal{D}_{*, \mathbf{a}}^{[m ; K]}\left(\mathbb{R}^{n}\right)$ is given by the seminorms
$\left\{\|\quad\|_{q, s}\right\}_{q>m, s \geq 0}$ defined as

$$
\begin{equation*}
\|\phi\|_{q, s}=\sup _{\mathbf{x}-\mathbf{a} \in K} \sup _{|\mathbf{p}| \leq s} \frac{\left|\frac{\partial^{\mathbf{p}} \phi}{\partial \mathbf{x}}(\mathbf{a}+\mathbf{x})-\sum_{j=m-|\mathbf{p}|}^{q-1} a_{j, \mathbf{p}}(\mathbf{w}) r^{j}\right|}{r^{q}}, \tag{2.2}
\end{equation*}
$$

where $\mathbf{x}=r \mathbf{w}$ and

$$
\begin{equation*}
\frac{\partial^{\mathbf{p}} \phi}{\partial \mathbf{x}}(\mathbf{a}+\mathbf{x}) \sim \sum_{j=m-|\mathbf{p}|}^{\infty} a_{j, \mathbf{p}}(\mathbf{w}) r^{j} \tag{2.3}
\end{equation*}
$$

The topology of $\mathcal{D}_{*, \mathbf{a}}^{[m]}\left(\mathbb{R}^{n}\right)$ is the inductive limit topology of the $\mathcal{D}_{*, \mathbf{a}}^{[m ; K]}\left(\mathbb{R}^{n}\right)$ as $K \nearrow \infty$. The topology of $\mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right)$ is the inductive limit topology of the $\mathcal{D}_{*, \mathbf{a}}^{[m]}\left(\mathbb{R}^{n}\right)$ as $m \searrow-\infty$.

A sequence $\left\{\phi_{l}\right\}_{l=0}^{\infty}$ in $\mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right)$ converges to $\psi$ if and only there exists $l_{0} \geq 0$, an integer $m$, and a compact set $K$ with a in its interior, such that $\phi_{l} \in \mathcal{D}_{*, \mathbf{a}}^{[m ; K]}\left(\mathbb{R}^{n}\right)$ for $l \geq l_{0}$ and $\left\|\psi-\phi_{l}\right\|_{q, s} \rightarrow 0$ as $l \rightarrow \infty$ if $q>m, s \geq 0$. Notice that if $\left\{\phi_{l}\right\}_{l=0}^{\infty}$ converges to $\psi$ in $\mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right)$ then $\phi_{l}$ and the corresponding derivatives converge uniformly to $\psi$ and its derivatives in any set of the form $\mathbb{R}^{n} \backslash B$, where $B$ is a ball with center at $\mathbf{a}$; in fact, $r^{|\mathbf{p}|-m}(\partial / \partial \mathbf{x})^{\mathbf{p}} \phi_{l}$ converges uniformly to $r^{|\mathbf{p}|-m}(\partial / \partial \mathbf{x})^{\mathbf{p}} \psi$ over all $\mathbb{R}^{n}$. Furthermore, if $\left\{a_{j}^{l}\right\}$ are the coefficients of the expansion of $\phi_{l}$ and $\left\{b_{j}\right\}$ are those for $\psi$, then $a_{j}^{l} \rightarrow b_{j}$ in the space $\mathcal{D}(\mathbb{S})$ for each $j \geq m$.

We can now consider distributions in a space with one thick point, the "thick distributions."

Definition 2.3. The space of distributions on $\mathbb{R}^{n}$ with a thick point at $\mathbf{x}=\mathbf{a}$ is the dual space of $\mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right)$. We denote it $\mathcal{D}_{*, \mathbf{a}}^{\prime}\left(\mathbb{R}^{n}\right)$, or just as $\mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$ when $\mathbf{a}=0$.

Observe that $\mathcal{D}\left(\mathbb{R}^{n}\right)$, the space of standard test functions, is a closed subspace of $\mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right)$; we denote by

$$
\begin{equation*}
i: \mathcal{D}\left(\mathbb{R}^{n}\right) \rightarrow \mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right), \tag{2.4}
\end{equation*}
$$

the inclusion map and by

$$
\begin{equation*}
\Pi: \mathcal{D}_{*, \mathbf{a}}^{\prime}\left(\mathbb{R}^{n}\right) \rightarrow \mathcal{D}^{\prime}\left(\mathbb{R}^{n}\right) \tag{2.5}
\end{equation*}
$$

the projection operator, dual of the inclusion (2.4).
The derivatives of thick distributions are defined in much the same way as the usual distributional derivatives, that is, by duality.

Definition 2.4. If $f \in \mathcal{D}_{*, \mathbf{a}}^{\prime}\left(\mathbb{R}^{n}\right)$ then its thick distributional derivative $\partial^{*} f / \partial x_{j}$ is defined as

$$
\begin{equation*}
\left\langle\frac{\partial^{*} f}{\partial x_{j}}, \phi\right\rangle=-\left\langle f, \frac{\partial \phi}{\partial x_{j}}\right\rangle, \quad \phi \in \mathcal{D}_{*, \mathbf{a}}\left(\mathbb{R}^{n}\right) \tag{2.6}
\end{equation*}
$$

We denote by $\mathcal{E}_{*}\left(\mathbb{R}^{n}\right)$ the space of smooth functions in $\mathbb{R}^{n} \backslash\{\mathbf{a}\}$ that have a strong asymptotic expansion of the form 2.1); alternatively, $\psi \in \mathcal{\mathcal { E } _ { * }}\left(\mathbb{R}^{n}\right)$ if $\psi=\psi_{1}+\psi_{2}$, where $\psi_{1} \in \mathcal{E}\left(\mathbb{R}^{n}\right)$, the space of all smooth functions in $\mathbb{R}^{n}$, and where $\psi_{2} \in \mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$. The space $\mathcal{E}_{*}\left(\mathbb{R}^{n}\right)$ is the space of multipliers of $\mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$ and of $\mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$. Furthermore [16], the product rule for derivatives holds,

$$
\begin{equation*}
\frac{\partial^{*}(\psi f)}{\partial x_{j}}=\frac{\partial \psi}{\partial x_{j}} f+\psi \frac{\partial^{*} f}{\partial x_{j}} \tag{2.7}
\end{equation*}
$$

if $f$ is a thick distribution and $\psi$ is a multiplier. Notice that $\partial \psi / \partial x_{j}$ is the ordinary derivative in 2.7).

Let $g(\mathbf{w})$ is a distribution in $\mathbb{S}$. The thick delta function of degree $q$, denoted as $g \delta_{*}^{[q]}$, or as $g(\mathbf{w}) \delta_{*}^{[q]}$, acts on a thick test function $\phi(\mathbf{x})$ as

$$
\begin{equation*}
\left\langle g \delta_{*}^{[q]}, \phi\right\rangle_{\mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right) \times \mathcal{D}_{*}\left(\mathbb{R}^{n}\right)}=\frac{1}{C}\left\langle g(\mathbf{w}), a_{q}(\mathbf{w})\right\rangle_{\mathcal{D}^{\prime}(\mathbb{S}) \times \mathcal{D}(\mathbb{S})}, \tag{2.8}
\end{equation*}
$$

where $\phi(r \mathbf{w}) \sim \sum_{j=m}^{\infty} a_{j}(\mathbf{w}) r^{j}$, as $r \rightarrow 0^{+}$, and where

$$
\begin{equation*}
C=\frac{2 \pi^{n / 2}}{\Gamma(n / 2)} \tag{2.9}
\end{equation*}
$$

is the surface area of the unit sphere $\mathbb{S}$ of $\mathbb{R}^{n}$. If $g$ is locally integrable function in $\mathbb{S}$, then

$$
\begin{equation*}
\left\langle g \delta_{*}^{[q]}, \phi\right\rangle_{\mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right) \times \mathcal{D}_{*}\left(\mathbb{R}^{n}\right)}=\frac{1}{C} \int_{\mathbb{S}} g(\mathbf{w}) a_{q}(\mathbf{w}) \mathrm{d} \sigma(\mathbf{w}) . \tag{2.10}
\end{equation*}
$$

Thick deltas of order 0 are called just thick deltas, and we shall use the notation $g \delta_{*}$ instead of $g \delta_{*}^{[0]}$.

Let $g \in \mathcal{D}^{\prime}(\mathbb{S})$. Then

$$
\begin{equation*}
\frac{\partial^{*}}{\partial x_{j}}\left(g \delta_{*}^{[q]}\right)=\left(\frac{\delta g}{\delta x_{j}}-(q+n) n_{j} g\right) \delta_{*}^{[q+1]} \tag{2.11}
\end{equation*}
$$

Here $\delta g / \delta x_{j}$ is the $\delta$-derivative of $g$ 4, 6; in general the $\delta$-derivatives can be applied to functions and distributions defined only on a smooth hypersurface $\Sigma$ of $\mathbb{R}^{n}$. Suppose now that the surface is $\mathbb{S}$, the unit sphere in $\mathbb{R}^{n}$ and let $f$ be a smooth function defined in $\mathbb{S}$, that is, $f(\mathbf{w})$ is defined if $\mathbf{w} \in \mathbb{R}^{n}$ satisfies $|\mathbf{w}|=1$. Observe that the expressions $\partial f / \partial x_{j}$ are not defined and, likewise, if $\mathbf{w}=\left(w_{j}\right)_{1 \leq j \leq n}$ the expressions $\partial f / \partial w_{j}$ do not make sense either; the derivatives that are always defined and that one should consider are the
$\delta f / \delta x_{j}, 1 \leq j \leq n$. Let $F_{0}$ be the extension of $f$ to $\mathbb{R}^{n} \backslash\{\mathbf{0}\}$ that is homogeneous of degree 0 , namely, $F_{0}(\mathbf{x})=f(\mathbf{x} / r)$ where $r=|\mathbf{x}|$. Then 16]

$$
\begin{equation*}
\frac{\delta f}{\delta x_{j}}=\left.\frac{\partial F_{0}}{\partial x_{j}}\right|_{\mathbb{S}} \tag{2.12}
\end{equation*}
$$

Also, if we use polar coordinates, $\mathbf{x}=r \mathbf{w}$, so that $F_{0}(\mathbf{x})=f(\mathbf{w})$, then $\partial F_{0} / \partial x_{j}$ is homogeneous of degree -1 , and actually $\partial F_{0} / \partial x_{j}=r^{-1} \delta f / \delta x_{j}$ if $\mathbf{x} \neq \mathbf{0}$.

The matrix $\mu=\left(\mu_{i j}\right)_{1 \leq i, j \leq n}$, where $\mu_{i j}=\delta n_{i} / \delta x_{j}$, plays an important role in the study of distributions on a surface $\Sigma$. If $\Sigma=\mathbb{S}$ then $\mu_{i j}=\delta n_{i} / \delta x_{j}=$ $\delta_{i j}-n_{i} n_{j}$. Observe that $\mu_{i j}=\mu_{j i}$, an identity that holds in any surface.

The differential operators $\delta f / \delta x_{j}$ are initially defined if $f$ is a smooth function defined on $\Sigma$, but we can also define them when $f$ is a distribution. We can do this if we use the fact that smooth functions are dense in the space of distributions on $\Sigma$.

## 3. The thick distribution $\mathcal{P} f(1)$

Let us consider one of the simplest functions, namely, the function 1, defined in $\mathbb{R}^{n}$. Naturally this function is locally integrable, and thus it defines a regular distribution, also denoted as 1 , and the ordinary derivatives and the distributional derivatives both coincide and give the value 0 . On the other hand, 1 does not automatically give an element of $\mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$ since if $\phi \in \mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$ the integral $\int_{\mathbb{R}^{n}} \phi(\mathbf{x}) \mathrm{d} \mathbf{x}$ could be divergent, and thus we consider the spherical finite part thick distribution $\mathcal{P} f(1)$ given as

$$
\begin{equation*}
\langle\mathcal{P} f(1), \phi\rangle=\text { F.p. } \int_{\mathbb{R}^{n}} \phi(\mathbf{x}) \mathrm{d} \mathbf{x}=\text { F.p. } \lim _{\varepsilon \rightarrow 0^{+}} \int_{|\mathbf{x}| \geq \varepsilon} \phi(\mathbf{x}) \mathrm{d} \mathbf{x} . \tag{3.1}
\end{equation*}
$$

The derivatives of $\mathcal{P} f(1)$ do not vanish, since actually we have the following formula 16 .

Lemma 3.1. In $\mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$,

$$
\begin{equation*}
\frac{\partial^{*}}{\partial x_{i}}(\mathcal{P} f(1))=C n_{i} \delta_{*}^{[-n+1]}, \tag{3.2}
\end{equation*}
$$

where $C$ is given by 2.9.
Proof. One can find a proof of a more general statement in [16, but in this simpler case the proof can be written as follows,

$$
\begin{aligned}
\left\langle\frac{\partial^{*}}{\partial x_{i}}(\mathcal{P} f(1)), \phi\right\rangle & =-\left\langle\mathcal{P} f(1), \frac{\partial \phi}{\partial x_{i}}\right\rangle \\
& =- \text { F.p. } \lim _{\varepsilon \rightarrow 0^{+}} \int_{|\mathbf{x}| \geq \varepsilon} \frac{\partial \phi}{\partial x_{i}} \mathrm{~d} \mathbf{x} \\
& =\text { F.p. } \lim _{\varepsilon \rightarrow 0^{+}} \int_{\varepsilon \mathbb{S}^{n}-1} n_{i} \phi \mathrm{~d} \sigma
\end{aligned}
$$

so that if $\phi \in \mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$ has the expansion $\phi(\mathbf{x}) \sim \sum_{j=m}^{\infty} a_{j}(\mathbf{w}) r^{j}$, as $\mathbf{x} \rightarrow \mathbf{0}$, then

$$
\int_{\varepsilon \mathbb{S}^{n-1}} n_{i} \phi \mathrm{~d} \sigma \sim \sum_{j=m}^{\infty}\left(\int_{\mathbb{S}} n_{i} a_{j}(\mathbf{w}) \mathrm{d} \sigma(\mathbf{w})\right) \varepsilon^{n-1+j}
$$

as $\varepsilon \rightarrow 0^{+}$. The finite part of the limit is equal to the coefficient of $\varepsilon^{0}$, thus

$$
\text { F.p. } \begin{aligned}
\lim _{\varepsilon \rightarrow 0} \int_{\varepsilon \mathbb{S}^{n-1}} n_{i} \phi \mathrm{~d} \sigma & =\int_{\mathbb{S}} n_{i} a_{1-n}(\mathbf{w}) \mathrm{d} \sigma(\mathbf{w}) \\
& =\left\langle C n_{i} \delta_{*}^{[1-n]}, \phi\right\rangle
\end{aligned}
$$

as required.
If $\psi \in \mathcal{E}_{*}\left(\mathbb{R}^{n}\right)$ is a multiplier of $\mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$, then we define, in a similar way, the thick distribution $\mathcal{P} f(\psi) \in \mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$, and we clearly have the useful formula

$$
\begin{equation*}
\mathcal{P} f(\psi)=\psi \mathcal{P} f(1), \tag{3.3}
\end{equation*}
$$

which immediately gives the thick distributional derivative of $\mathcal{P} f(\psi)$ as

$$
\frac{\partial^{*}}{\partial x_{i}}(\mathcal{P} f(\psi))=\frac{\partial \psi}{\partial x_{i}} \mathcal{P} f(1)+\psi \frac{\partial^{*}}{\partial x_{i}}(\mathcal{P} f(1))
$$

so that we obtain the ensuing formula.
Proposition 3.2. If $\psi \in \mathcal{E}_{*}\left(\mathbb{R}^{n}\right)$ then

$$
\begin{equation*}
\frac{\partial^{*}}{\partial x_{i}}(\mathcal{P} f(\psi))=\mathcal{P} f\left(\frac{\partial \psi}{\partial x_{i}}\right)+C n_{i} \psi \delta_{*}^{[1-n]} . \tag{3.4}
\end{equation*}
$$

Notice that, in general, the term $C n_{i} \psi \delta_{*}^{[1-n]}$ is not a thick delta of order $1-n$. Indeed, let us now consider the case when $\psi \in \mathcal{E}_{*}\left(\mathbb{R}^{n}\right)$ is homogeneous of order $k \in \mathbb{Z}$. Then $\psi(\mathbf{x})=r^{k} \psi_{0}(\mathbf{x})$, where $\psi_{0}$ is homogeneous of order 0 . Since $r^{k} \delta_{*}^{[q]}=\delta_{*}^{[q-k]}$ [16, Eqn. (5.16)] we obtain the following particular case of (3.4), where now the term $C n_{i} \psi_{0} \delta_{*}^{[1-n-k]}$ is a thick delta of order $1-n-k$.

Proposition 3.3. If $\psi \in \mathcal{\mathcal { E } _ { * }}\left(\mathbb{R}^{n}\right)$ is homogeneous of order $k \in \mathbb{Z}$, then

$$
\begin{equation*}
\frac{\partial^{*}}{\partial x_{i}}(\mathcal{P} f(\psi))=\mathcal{P} f\left(\frac{\partial \psi}{\partial x_{i}}\right)+C n_{i} \psi_{0} \delta_{*}^{[1-n-k]} \tag{3.5}
\end{equation*}
$$

where $\psi_{0}(\mathbf{x})=|\mathbf{x}|^{-k} \psi(\mathbf{x})$.
If we now apply the projection $\Pi$ onto the usual distribution space $\mathcal{D}^{\prime}\left(\mathbb{R}^{n}\right)$, we obtain the formula for the distributional derivatives of homogeneous distributions. Observe first that if $k>-n$ then $\psi$ is integrable at the origin, and thus $\psi$ is a regular distribution and $\Pi(\mathcal{P} f(\psi))=\psi$. If $k \leq-n$ then $\Pi(\mathcal{P} f(\psi))=\mathcal{P} f(\psi)$, since in that case the integral $\int_{\mathbb{R}^{n}} \psi(\mathbf{x}) \phi(\mathbf{x}) \mathrm{d} \mathbf{x}$ would
be divergent, in general, if $\phi \in \mathcal{D}\left(\mathbb{R}^{n}\right)$. A particularly interesting case is when $k=-n$, since if $\psi$ is homogeneous of degree $-n$ and

$$
\begin{equation*}
\int_{\mathbb{S}} \psi(\mathbf{w}) \mathrm{d} \sigma(\mathbf{w})=0 \tag{3.6}
\end{equation*}
$$

then the principal value of the integral

$$
\begin{equation*}
\text { p.v. } \int_{\mathbb{R}^{n}} \psi(\mathbf{x}) \phi(\mathbf{x}) \mathrm{d} \mathbf{x}=\lim _{\varepsilon \rightarrow 0^{+}} \int_{|\mathbf{x}| \geq \varepsilon} \psi(\mathbf{x}) \phi(\mathbf{x}) \mathrm{d} \mathbf{x} \tag{3.7}
\end{equation*}
$$

actually exists for each $\phi \in \mathcal{D}\left(\mathbb{R}^{n}\right)$, so that $\mathcal{P} f(\psi)=$ p.v. $(\psi)$, the principal value distribution. Note however that if $\Sigma$ is a closed surface in $\mathbb{R}^{n}$ that encloses the origin, described by an equation of the form $g(\mathbf{x})=1$, where $g(\mathbf{x})$ is continuous in $\mathbb{R}^{n} \backslash\{\mathbf{0}\}$ and homogeneous of degree 1 , then $\left\langle\mathcal{R}_{\Sigma}(\psi(\mathbf{x})), \phi(\mathbf{x})\right\rangle=\lim _{\varepsilon \rightarrow 0} \int_{g(\mathbf{x}) \geqslant \varepsilon} \psi(\mathbf{x}) \phi(\mathbf{x}) \mathrm{d} \mathbf{x}$, defines another regularization of $\psi$, but in general $\mathcal{R}_{\Sigma}(\psi(\mathbf{x})) \neq$ p.v. $(\psi(\mathbf{x}))$ [15], a fact observed by Farassat [8, who indicated its importance in numerical computations, and studied by several authors [11, 15].

Condition (3.6) holds whenever $\psi=\partial \xi / \partial x_{j}$ for some $\xi$ homogeneous of order $-n+1$.

Proposition 3.4. Let $\psi$ be homogeneous of order $k \in \mathbb{Z}$ in $\mathbb{R}^{n} \backslash\{\mathbf{0}\}$. Then, in $\mathcal{D}^{\prime}\left(\mathbb{R}^{n}\right)$ the distributional derivative $\bar{\partial} \psi / \partial x_{i}$ is given as follows:

$$
\begin{equation*}
\frac{\bar{\partial} \psi}{\partial x_{i}}=\frac{\partial \psi}{\partial x_{i}}, \quad k>1-n \tag{3.8}
\end{equation*}
$$

equality of regular distributions;

$$
\begin{equation*}
\frac{\bar{\partial} \psi}{\partial x_{i}}=\text { p.v. }\left(\frac{\partial \psi}{\partial x_{i}}\right)+A \delta(\mathbf{x}), \quad k=1-n \tag{3.9}
\end{equation*}
$$

where $A=\int_{\mathbb{S}} n_{i} \psi_{0}(\mathbf{w}) \mathrm{d} \sigma(\mathbf{w})=\left\langle\psi_{0}, n_{i}\right\rangle_{\mathcal{D}^{\prime}(\mathbb{S}) \times \mathcal{D}(\mathbb{S})}$, while

$$
\begin{equation*}
\frac{\bar{\partial} \psi}{\partial x_{i}}=\mathcal{P} f\left(\frac{\partial \psi}{\partial x_{i}}\right)+D(\mathbf{x}), \quad k<1-n \tag{3.10}
\end{equation*}
$$

where $D(\mathbf{x})$ is a homogeneous distribution of order $k-1$ concentrated at the origin and given by

$$
\begin{equation*}
D(\mathbf{x})=(-1)^{-k-n+1} \sum_{j_{1}+\cdots+j_{n}=-k-n+1} \frac{\left\langle n_{i} \psi_{0}, \mathbf{w}^{\left(j_{1}, \ldots, j_{n}\right)}\right\rangle}{j_{1}!\cdots j_{n}!} \mathbf{D}^{\left(j_{1}, \ldots, j_{n}\right)} \delta(\mathbf{x}) \tag{3.11}
\end{equation*}
$$

Proof. It follows from (3.4) if we observe [16, Prop. 4.7] that if $g \in \mathcal{D}^{\prime}(\mathbb{S})$ then

$$
\begin{equation*}
\Pi\left(g \delta_{*}^{[q]}\right)=\frac{(-1)^{q}}{C} \sum_{j_{1}+\cdots+j_{n}=q} \frac{\left\langle g(\mathbf{w}), \mathbf{w}^{\left(j_{1}, \ldots, j_{n}\right)}\right\rangle}{j_{1}!\cdots j_{n}!} \mathbf{D}^{\left(j_{1}, \ldots, j_{n}\right)} \delta(\mathbf{x}) \tag{3.12}
\end{equation*}
$$

and, in particular,

$$
\begin{equation*}
\Pi\left(g \delta_{*}\right)=\frac{1}{C}\langle g(\mathbf{w}), 1\rangle \delta(\mathbf{x}), \tag{3.13}
\end{equation*}
$$

if $q=0$.
Our next task is to compute the second order thick derivatives of homogeneous distributions. Indeed, if $\psi$ is homogeneous of degree $k$ then we can iterate the formula (3.5) to obtain

$$
\begin{align*}
\frac{\partial^{* 2}}{\partial x_{i} \partial x_{j}} & (\mathcal{P} f(\psi))=\frac{\partial^{*}}{\partial x_{i}}\left(\mathcal{P} f\left(\frac{\partial \psi}{\partial x_{j}}\right)+C n_{j} \psi_{0} \delta_{*}^{[1-n-k]}\right)  \tag{3.14}\\
& =\mathcal{P} f\left(\frac{\partial^{2} \psi}{\partial x_{i} \partial x_{j}}\right)+C n_{i} \xi_{0} \delta_{*}^{[2-n-k]}+\frac{\partial^{*}}{\partial x_{i}}\left(C n_{j} \psi_{0} \delta_{*}^{[1-n-k]}\right)
\end{align*}
$$

where $\xi=\partial \psi / \partial x_{j}$ is homogeneous of degree $k-1$ and $\xi_{0}(\mathbf{x})=|\mathbf{x}|^{1-k} \xi(\mathbf{x})$ is the associated function which is homogeneous of degree 0 . Use of 2.11) allows us to write

$$
\begin{align*}
\frac{\partial^{*}}{\partial x_{i}} & \left(C n_{j} \psi_{0} \delta_{*}^{[1-n-k]}\right)=C\left(\frac{\delta}{\delta x_{i}}\left(n_{j} \psi_{0}\right)+(k-1) n_{i} n_{j} \psi_{0}\right) \delta_{*}^{[2-n-k]}  \tag{3.15}\\
& =C\left(\left(\delta_{i j}-n_{i} n_{j}\right) \psi_{0}+n_{j} \frac{\delta \psi_{0}}{\delta x_{i}}+(k-1) n_{i} n_{j} \psi_{0}\right) \delta_{*}^{[2-n-k]} \\
& =C\left(\left(\delta_{i j}+(k-2) n_{i} n_{j}\right) \psi_{0}+n_{j} \frac{\delta \psi_{0}}{\delta x_{i}}\right) \delta_{*}^{[2-n-k]}
\end{align*}
$$

while the equation $\psi=r^{k} \psi_{0}$ yields $\partial \psi / \partial x_{j}=r^{k-1}\left\{k n_{j} \psi_{0}+\delta \psi_{0} / \delta x_{j}\right\}$, so that

$$
\begin{equation*}
\xi_{0}=k n_{j} \psi_{0}+\frac{\delta \psi_{0}}{\delta x_{j}} \tag{3.16}
\end{equation*}
$$

Collecting terms we thus obtain the following formula.
Proposition 3.5. If $\psi \in \mathcal{\mathcal { E } _ { * }}\left(\mathbb{R}^{n}\right)$ is homogeneous of order $k \in \mathbb{Z}$, then

$$
\begin{align*}
& \frac{\partial^{* 2}}{\partial x_{i} \partial x_{j}}(\mathcal{P} f(\psi))=\mathcal{P} f\left(\frac{\partial^{2} \psi}{\partial x_{i} \partial x_{j}}\right)  \tag{3.17}\\
& \quad+C\left(\left(\delta_{i j}+2(k-1) n_{i} n_{j}\right) \psi_{0}+n_{j} \frac{\delta \psi_{0}}{\delta x_{i}}+n_{i} \frac{\delta \psi_{0}}{\delta x_{j}}\right) \delta_{*}^{[2-n-k]}
\end{align*}
$$

where $\psi_{0}(\mathbf{x})=|\mathbf{x}|^{-k} \psi(\mathbf{x})$.
Projection onto $\mathcal{D}^{\prime}\left(\mathbb{R}^{n}\right)$ of 3.17) gives the formula for the distributional derivatives $\bar{\partial}^{2} / \partial x_{i} \partial x_{j}(\mathcal{P} f(\psi))$ if $\psi \in \mathcal{E}_{*}\left(\mathbb{R}^{n}\right)$ is homogeneous of order $k \in \mathbb{Z}$. In case $k=2-n$ we obtain the following formula.

Proposition 3.6. If $\psi \in \mathcal{E}_{*}\left(\mathbb{R}^{n}\right)$ is homogeneous of order $2-n$, then

$$
\begin{equation*}
\frac{\bar{\partial}^{2}}{\partial x_{i} \partial x_{j}}(\psi)=\text { p.v. }\left(\frac{\partial^{2} \psi}{\partial x_{i} \partial x_{j}}\right)+B \delta(\mathbf{x}) \tag{3.18}
\end{equation*}
$$

where

$$
\begin{equation*}
B=\left\langle\psi_{0}, 2 n_{i} n_{j}-\delta_{i j}\right\rangle_{\mathcal{D}^{\prime}(\mathbb{S}) \times \mathcal{D}(\mathbb{S})} \tag{3.19}
\end{equation*}
$$

Proof. If we apply the operator $\Pi$ to 3.17 and employ 3.13 we obtain 3.18 with

$$
B=\left\langle\left(\delta_{i j}+2(k-1) n_{i} n_{j}\right) \psi_{0}+n_{j} \frac{\delta \psi_{0}}{\delta x_{j}}+n_{i} \frac{\delta \psi_{0}}{\delta x_{j}}, 1\right\rangle_{\mathcal{D}^{\prime}(\mathbb{S}) \times \mathcal{D}(\mathbb{S})}
$$

But [16, (2.6)] yields

$$
\begin{equation*}
\left\langle n_{j} \frac{\delta \psi_{0}}{\delta x_{j}}, 1\right\rangle_{\mathcal{D}^{\prime}(\mathbb{S}) \times \mathcal{D}(\mathbb{S})}=\left\langle\psi_{0}, n n_{i} n_{j}-\delta_{i j}\right\rangle_{\mathcal{D}^{\prime}(\mathbb{S}) \times \mathcal{D}(\mathbb{S})} \tag{3.20}
\end{equation*}
$$

and (3.19) follows since $k=2-n$.
We would like to observe that while $\psi_{0}$ has been supposed smooth, a continuity argument immediately gives that $\psi_{0}$ could be any distribution of $\mathcal{D}^{\prime}\left(\mathbb{R}^{n} \backslash\{\mathbf{0}\}\right)$ that is homogeneous of degree 0.

## 4. Bowen's formula

If we apply formula $\sqrt{3.9}$ to the function $\psi=n_{j_{1}} \cdots n_{j_{k}} / r^{2}$, which is homogeneous of degree -2 in $\mathbb{R}^{3}$ we obtain at once that

$$
\begin{align*}
& \frac{\bar{\partial}}{\partial x_{i}}\left(\frac{n_{j_{1}} \cdots n_{j_{k}}}{r^{2}}\right)=  \tag{4.1}\\
& \text { p.v. }\left(\left\{\sum_{q=1}^{k} \delta_{i j_{q}} \frac{n_{j_{1}} \cdots n_{j_{k}}}{n_{j_{q}}}-(k+2) n_{i} n_{j_{1}} \cdots n_{j_{k}}\right\} \frac{1}{r^{3}}\right)+A \delta(\mathbf{x})
\end{align*}
$$

where

$$
\begin{equation*}
A=\int_{\mathbb{S}} n_{i} n_{j_{1}} \cdots n_{j_{k}} \mathrm{~d} \sigma(\mathbf{w}) \tag{4.2}
\end{equation*}
$$

This integral was computed in [5, (3.13)], the result being

$$
\begin{equation*}
A=\frac{2 \Gamma((a+1) / 2) \Gamma((b+1) / 2) \Gamma((c+1) / 2)}{\Gamma((a+b+c+3) / 2)} \tag{4.3}
\end{equation*}
$$

if $n_{i} n_{j_{1}} \cdots n_{j_{k}}=n_{1}^{a} n_{2}^{b} n_{3}^{c}$, and $a, b$, or $c$ are even, while $A=0$ if any exponent is odd. Bowen [2, Eqn. (A5)] also computes the integral, and obtains a different but equivalent expression; in particular, his formula for $k=3$ reads as

$$
\begin{equation*}
A=\frac{4 \pi}{15}\left(\delta_{i j_{1}} \delta_{j_{2} j_{3}}+\delta_{i j_{2}} \delta_{j_{1} j_{3}}+\delta_{i j_{3}} \delta_{j_{1} j_{2}}\right) \tag{4.4}
\end{equation*}
$$

so that 4.3) or 4.4 would yield that if $(a, b, c)$ is a permutation of $(2,2,0)$ then $A=4 \pi / 15$ while if $(a, b, c)$ is a permutation of $(4,0,0)$ then $A=4 \pi / 5$.

Our main aim is to point out why the product rule for derivatives, as employed in [2] does not produce the correct result. Indeed, if we use [2, Eqn. (16)] written as

$$
\begin{equation*}
\frac{\bar{\partial}}{\partial x_{i}}\left(\frac{n_{j_{1}}}{r^{2}}\right)=\text { p.v. }\left(\frac{\delta_{i j_{1}}-3 n_{i} n_{j_{i}}}{r^{3}}\right)+\frac{4 \pi}{3} \delta_{i j_{1}} \delta(\mathbf{x}) \tag{4.5}
\end{equation*}
$$

and then try to proceed as in [2, Eqn. (18)],

$$
\begin{equation*}
\frac{\bar{\partial}}{\partial x_{i}}\left(\frac{n_{j_{1}} n_{j_{2}} n_{j_{3}}}{r^{2}}\right) " i=? " n_{j_{1}} n_{j_{2}} \frac{\bar{\partial}}{\partial x_{i}}\left(\frac{n_{j_{3}}}{r^{2}}\right)+\frac{n_{j_{3}}}{r^{2}} \frac{\bar{\partial}}{\partial x_{i}}\left(n_{j_{1}} n_{j_{2}}\right) . \tag{4.6}
\end{equation*}
$$

Thus 4.5 and the formula

$$
\begin{equation*}
\frac{\bar{\partial}}{\partial x_{i}}\left(n_{j_{1}} n_{j_{2}}\right)=\frac{\delta_{i j_{1}} n_{j_{2}}+\delta_{i j_{2}} n_{j_{1}}-2 n_{i} n_{j_{1}} n_{j_{2}}}{r} \tag{4.7}
\end{equation*}
$$

give

$$
\begin{equation*}
n_{j_{1}} n_{j_{2}} \frac{\bar{\partial}}{\partial x_{i}}\left(\frac{n_{j_{3}}}{r^{2}}\right)+\frac{n_{j_{3}}}{r^{2}} \frac{\bar{\partial}}{\partial x_{i}}\left(n_{j_{1}} n_{j_{2}}\right)=\text { "Normal" }+ \text { "Src" } \tag{4.8}
\end{equation*}
$$

where

$$
\begin{equation*}
" N o r m a l "=\text { p.v. }\left(\frac{\delta_{i j_{1}} n_{j_{2}} n_{j_{3}}+\delta_{i j_{2}} n_{j_{1}} n_{j_{3}}+\delta_{i j_{3}} n_{j_{1}} n_{j_{2}}-5 n_{i} n_{j_{1}} n_{j_{2}} n_{j_{3}}}{r^{3}}\right), \tag{4.9}
\end{equation*}
$$

coincides with the first term of 4.1 while

$$
\begin{equation*}
" S r c "=\frac{4 \pi}{3} \delta_{i j_{3}} n_{j_{1}} n_{j_{2}} \delta(\mathbf{x}) \tag{4.10}
\end{equation*}
$$

The right hand side of 4.10 is not a well defined distribution, of course, but Bowen suggested that we treat it as what we now call the projection of a thick distribution, that is, as

$$
\begin{equation*}
" \mathrm{Src} "=\Pi\left(\frac{4 \pi}{3} \delta_{i j_{3}} n_{j_{1}} n_{j_{2}} \delta_{*}\right)=\frac{4 \pi}{9} \delta_{i j_{3}} \delta_{j_{1} j_{2}} \delta(\mathbf{x}) \tag{4.11}
\end{equation*}
$$

since $\Pi\left(n_{j_{1}} n_{j_{2}} \delta_{*}\right)=(1 / 3) \delta_{j_{1} j_{2}} \delta(\mathbf{x})$ [16, Example 5.10]. In order to compare with (4.1) and (4.4) we observe that by symmetry the same result would be obtained if $j_{3}$ and $j_{1}$, or $j_{3}$ and $j_{2}$, are exchanged, so that if in the term "Src" we do these exchanges, add the results and divide by 3 , we would get

$$
\begin{equation*}
" S r c S y m "=\frac{4 \pi}{27}\left(\delta_{i j_{1}} \delta_{j_{2} j_{3}}+\delta_{i j_{2}} \delta_{j_{1} j_{3}}+\delta_{i j_{3}} \delta_{j_{1} j_{2}}\right) \delta(\mathbf{x}), \tag{4.12}
\end{equation*}
$$

and thus the symmetric version of the 4.8 is "Normal" + "SrcSym", which of course is different from 4.1) since the coefficient in 4.4 is $4 \pi / 15$, while that
in (4.12) is $4 \pi / 27$. Therefore, the relation " $i=$ ?" in 4.6 cannot be replaced by $=$.

Hence the product rule for derivatives fails in this case. The question is why? Indeed, when computing the right side of 4.6), that is, the left side of 4.8), we found just one irregular product, namely $n_{j_{1}} n_{j_{2}} \delta(\mathbf{x})$, but using the average value $(1 / 3) \delta_{j_{1} j_{2}} \delta(\mathbf{x})$ seems quite reasonable.

In order to see what went wrong let us compute $\bar{\partial} / \partial x_{i}\left(n_{j_{1}} n_{j_{2}} n_{j_{3}} / r^{2}\right)$ by computing the thick derivative $\partial^{*} / \partial x_{i} \mathcal{P} f\left(n_{j_{1}} n_{j_{2}} n_{j_{3}} / r^{2}\right)$, applying the product rule for thick derivatives, and then taking the projection $\pi$ of this. We have,

$$
\begin{aligned}
\frac{\partial^{*}}{\partial x_{i}} \mathcal{P} f\left(\frac{n_{j_{1}} n_{j_{2}} n_{j_{3}}}{r^{2}}\right) & =\frac{\partial^{*}}{\partial x_{i}}\left[n_{j_{1}} n_{j_{2}} \mathcal{P} f\left(\frac{n_{j_{3}}}{r^{2}}\right)\right] \\
& =n_{j_{1}} n_{j_{2}} \frac{\partial^{*}}{\partial x_{i}} \mathcal{P} f\left(\frac{n_{j_{3}}}{r^{2}}\right)+\frac{\partial\left(n_{j_{1}} n_{j_{2}}\right)}{\partial x_{i}} \mathcal{P} f\left(\frac{n_{j_{3}}}{r^{2}}\right),
\end{aligned}
$$

and taking 3.5 into account, we obtain

$$
\begin{aligned}
& n_{j_{1}} n_{j_{2}}\left\{\mathcal{P} f\left(\frac{\delta_{i j_{3}}-3 n_{i} n_{j_{3}}}{r^{3}}\right)+4 \pi n_{j_{3}} n_{i} \delta_{*}\right\} \\
& +\frac{\delta_{i j_{1}} n_{j_{2}}+\delta_{i j_{2}} n_{j_{1}}-2 n_{i} n_{j_{1}} n_{j_{2}}}{r} \mathcal{P} f\left(\frac{n_{j_{3}}}{r^{2}}\right)
\end{aligned}
$$

that is, $\partial^{*} / \partial x_{i} \mathcal{P} f\left(n_{j_{1}} n_{j_{2}} n_{j_{3}} / r^{2}\right)$ equals

$$
\begin{equation*}
\mathcal{P f}\left(\frac{\delta_{i j_{1}} n_{j_{2}} n_{j_{3}}+\delta_{i j_{2}} n_{j_{1}} n_{j_{3}}+\delta_{i j_{3}} n_{j_{1}} n_{j_{2}}-5 n_{i} n_{j_{1}} n_{j_{2}} n_{j_{3}}}{r^{3}}\right) \tag{4.13}
\end{equation*}
$$

Applying the projection operator $\Pi$ we obtain that the $\mathcal{P} f$ becomes a p.v., so that the term "Normal" given by (4.9) is obtained, while (3.13) yields that the projection of thick delta is exactly $A \delta(\mathbf{x})$ where $A=\int_{\mathbb{S}} n_{i} n_{j_{1}} n_{j_{2}} n_{j_{3}} \mathrm{~d} \sigma(\mathbf{w})$, that is, the correct term

$$
\frac{4 \pi}{15}\left(\delta_{i j_{1}} \delta_{j_{2} j_{3}}+\delta_{i j_{2}} \delta_{j_{1} j_{3}}+\delta_{i j_{3}} \delta_{j_{1} j_{2}}\right) \delta(\mathbf{x}) .
$$

The reason we now obtain the correct result is while it is true that $\Pi\left(n_{j_{1}} n_{j_{2}} \delta_{*}\right)=$ $(1 / 3) \delta_{j_{1} j_{2}} \delta(\mathbf{x})$ and that $\Pi\left(n_{j_{3}} n_{i} \delta_{*}\right)=(1 / 3) \delta_{i j_{3}} \delta(\mathbf{x})$, it is not true that the projection $\Pi\left(4 \pi n_{j_{1}} n_{j_{2}} n_{j_{3}} n_{i} \delta_{*}\right)$ can be obtained as $4 \pi(1 / 3) \delta_{i j_{3}} \Pi\left(n_{j_{1}} n_{j_{2}} \delta_{*}\right)$ nor as $4 \pi(1 / 3) \delta_{j_{1} j_{2}} \Pi\left(n_{j_{3}} n_{i} \delta_{*}\right)$, and actually not even the symmetrization of such results, given by 4.12, works. Put in simple terms, it is not true that the average of a product is the product of the averages!

One can, alternatively, compute $\partial^{*} / \partial x_{i} \mathcal{P} f\left(n_{j_{1}} n_{j_{2}} n_{j_{3}} / r^{2}\right)$ as

$$
\begin{equation*}
\frac{\partial}{\partial x_{i}}\left(\frac{n_{j_{3}}}{r^{2}}\right) \mathcal{P} f\left(n_{j_{1}} n_{j_{2}}\right)+\left(\frac{n_{j_{3}}}{r^{2}}\right) \frac{\partial^{*}}{\partial x_{i}} \mathcal{P} f\left(n_{j_{1}} n_{j_{2}}\right) \tag{4.14}
\end{equation*}
$$

since

$$
\begin{equation*}
\frac{\partial^{*}}{\partial x_{i}} \mathcal{P} f\left(n_{j_{1}} n_{j_{2}}\right)=\mathcal{P} f\left(\frac{\delta_{i j_{1}} n_{j_{2}}+\delta_{i j_{2}} n_{j_{1}}-2 n_{i} n_{j_{1}} n_{j_{2}}}{r}\right)+4 \pi n_{j_{1}} n_{j_{2}} n_{i} \delta_{*}^{[-2]} \tag{4.15}
\end{equation*}
$$

Here the thick delta term in 4.14) is $4 \pi\left(n_{j_{3}} / r^{2}\right) n_{j_{1}} n_{j_{2}} n_{i} \delta_{*}^{[-2]}$, which becomes, as it should, $4 \pi n_{j_{1}} n_{j_{2}} n_{j_{3}} n_{i} \delta_{*}$.

Complications in the use of the product rule for derivatives in one variable were considered in [3] when analysing the formula [14]

$$
\begin{equation*}
\frac{d}{d x}\left(H^{n}(x)\right)=n H^{n-1}(x) \delta(x) \tag{4.16}
\end{equation*}
$$

where $H$ is the Heaviside function; see also [12].

## 5. Higher order derivatives

We now consider the computation of higher order derivatives in the space $\left(\mathcal{D}_{*}^{[0]}\left(\mathbb{R}^{n}\right)\right)^{\prime}$. If $f \in \mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$ then, of course, the thick derivative $\partial^{*} f / \partial x_{i}$ is defined by duality, that is,

$$
\begin{equation*}
\left\langle\frac{\partial^{*} f}{\partial x_{i}}, \phi\right\rangle=-\left\langle f, \frac{\partial \phi}{\partial x_{i}}\right\rangle \tag{5.1}
\end{equation*}
$$

for $\phi \in \mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$. Suppose now that $\mathcal{A}$ is a subspace of $\mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$ that has a topology such that the imbedding $i: \mathcal{A} \hookrightarrow \mathcal{D}_{*}\left(\mathbb{R}^{n}\right)$ is continuous; then the transpose $i^{T}: \mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right) \rightarrow \mathcal{A}^{\prime}$ is just the restriction operator $\Pi_{\mathcal{A}}$. If $\mathcal{A}$ is closed under the differentiation operators. Note that, the space $\mathcal{A}^{\prime}$ would be a space of (thick) distributions in the sense of Zemanian [18]. Then we can also define the derivative of any $f \in \mathcal{A}^{\prime}$, say $\partial_{\mathcal{A}} f / \partial x_{i}$, by employing (5.1) for $\phi \in \mathcal{A}$. Then

$$
\begin{equation*}
\Pi_{\mathcal{A}}\left(\frac{\partial^{*} f}{\partial x_{i}}\right)=\frac{\partial_{\mathcal{A}}}{\partial x_{i}}\left(\Pi_{\mathcal{A}}(f)\right), \tag{5.2}
\end{equation*}
$$

for any thick distribution $f \in \mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$. In the particular case when $\mathcal{A}=\mathcal{D}\left(\mathbb{R}^{n}\right)$, then $\partial_{\mathcal{A}} f / \partial x_{i}=\bar{\partial} f / \partial x_{i}$, the usual distributional derivative, and thus 5.2) becomes [16, Eqn. (5.22)],

$$
\begin{equation*}
\Pi\left(\frac{\partial^{*} f}{\partial x_{i}}\right)=\frac{\bar{\partial} \Pi(f)}{\partial x_{i}} \tag{5.3}
\end{equation*}
$$

What this means is that one can use thick distributional derivatives to compute $\partial_{\mathcal{A}} f / \partial x_{i}$, as we have already done to compute distributional derivatives.

When $\mathcal{A}$ is not closed under the differentiation operators then $\partial_{\mathcal{A}} f / \partial x_{i}$ cannot be defined by (5.1) if $f \in \mathcal{A}^{\prime}$ since in general $\partial \phi / \partial x_{i}$ does not belong to $\mathcal{A}$ and thus the right side of (5.1) is not defined. However, if $f \in \mathcal{A}^{\prime}$ has a canonical extension $\widetilde{f} \in \mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$ then we could define $\partial_{\mathcal{A}} f / \partial x_{i}$ as $\Pi_{\mathcal{A}}\left(\partial^{*} \widetilde{f} / \partial x_{i}\right)$. This applies, in particular when $\mathcal{A}=\mathcal{D}_{*}^{[0]}\left(\mathbb{R}^{n}\right)$ : if $f \in\left(\mathcal{D}_{*}^{[0]}\left(\mathbb{R}^{n}\right)\right)^{\prime}$ then $\partial_{0}^{*} f / \partial x_{i}=\partial_{\mathcal{A}} f / \partial x_{i}$ cannot be defined, in general, but if $f$ has a canonical extension $\widetilde{f} \in \mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$, then $\partial_{0}^{*} f / \partial x_{i}$ is understood as $\Pi_{\mathcal{D}_{*}^{[0]}\left(\mathbb{R}^{n}\right)}\left(\partial^{*} \widetilde{f} / \partial x_{i}\right)$.

Our aim is to point out that, in general, if $P=R S$ is the product of two differential operators with constant coefficients, then while, with obvious
notations, $P^{*}=R^{*} S^{*}, P_{\mathcal{A}}=R_{\mathcal{A}} S_{\mathcal{A}}$, if $\mathcal{A}$ is closed under differential operators, and $\bar{P}=\bar{R} \bar{S}$, it is not true that $P_{0}^{*}=R_{0}^{*} S_{0}^{*}$. Therefore the space $\left(\mathcal{D}_{*}^{[0]}\left(\mathbb{R}^{n}\right)\right)^{\prime}$ is not a convenient framework to generalize distributions to thick distributions; the whole $\mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{n}\right)$ is needed if we want a theory that includes the possibility of differentiation.

Example 5.1. Let us consider the second order derivatives of the distribution $\mathcal{P} f(1)$. Formula (3.17) yields

$$
\begin{equation*}
\frac{\partial^{* 2}}{\partial x_{i} \partial x_{j}}(\mathcal{P} f(1))=C\left(\delta_{i j}-2 n_{i} n_{j}\right) \delta_{*}^{[-n+2]} \tag{5.4}
\end{equation*}
$$

In particular, in $\mathbb{R}^{2}, \partial^{* 2} / \partial x_{i} \partial x_{j}(\mathcal{P} f(1))=2 \pi\left(\delta_{i j}-2 n_{i} n_{j}\right) \delta_{*}$. If we consider the function 1 as an element of $\left(\mathcal{D}_{*}^{[0]}\left(\mathbb{R}^{2}\right)\right)^{\prime}$ then it has the canonical extension $\mathcal{P} f(1) \in \mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{2}\right)$ and so

$$
\frac{\partial_{0}^{*}(1)}{\partial x_{j}}=\Pi_{\mathcal{D}_{*}^{[0]}\left(\mathbb{R}^{2}\right)}\left(2 \pi n_{j} \delta_{*}^{[-1]}\right)=0
$$

and consequently,

$$
\begin{equation*}
\frac{\partial_{0}^{*}}{\partial x_{i}}\left(\frac{\partial_{0}^{*}(1)}{\partial x_{j}}\right)=\frac{\partial_{0}^{*}}{\partial x_{i}}(0)=0 \neq 2 \pi\left(\delta_{i j}-2 n_{i} n_{j}\right) \delta_{*}=\frac{\partial^{* 2}(1)}{\partial x_{i} \partial x_{j}} . \tag{5.5}
\end{equation*}
$$

Observe that $\Pi\left(2 \pi\left(\delta_{i j}-2 n_{i} n_{j}\right) \delta_{*}\right)=0$, but observe also that this means very little.

Example 5.2. It was obtained in [16, Thm. 7.6] that in $\mathcal{D}_{*}^{\prime}\left(\mathbb{R}^{3}\right)$

$$
\begin{equation*}
\frac{\partial^{* 2} \mathcal{P} f\left(r^{-1}\right)}{\partial x_{i} \partial x_{j}}=\left(3 x_{i} x_{j}-\delta_{i j} r^{2}\right) \mathcal{P} f\left(r^{-5}\right)+4 \pi\left(\delta_{i j}-4 n_{i} n_{j}\right) \delta_{*} \tag{5.6}
\end{equation*}
$$

Since $\Pi\left(n_{i} n_{j} \delta_{*}\right)=(1 / 3) \delta_{i j} \delta(\mathbf{x})$ in $\mathbb{R}^{3}$, this yields the well known formula of Frahm (9]

$$
\begin{equation*}
\frac{\bar{\partial}^{2}}{\partial x_{i} \partial x_{j}}\left(\frac{1}{r}\right)=\text { p.v. }\left(\frac{3 x_{i} x_{j}-r^{2} \delta_{i j}}{r^{5}}\right)-\left(\frac{4 \pi}{3}\right) \delta_{i j} \delta(\mathbf{x}) \tag{5.7}
\end{equation*}
$$

We also immediately obtain that

$$
\begin{equation*}
\frac{\partial_{0}^{* 2} \mathcal{P} f\left(r^{-1}\right)}{\partial x_{i} \partial x_{j}}=\mathcal{P} f\left(\frac{3 x_{i} x_{j}-r^{2} \delta_{i j}}{r^{5}}\right)+4 \pi\left(\delta_{i j}-4 n_{i} n_{j}\right) \delta_{*} \tag{5.8}
\end{equation*}
$$

a formula that can also be proved by other methods 17. On the other hand, in [10] one can find the computation of

$$
\begin{equation*}
\frac{\partial_{0}^{*}}{\partial x_{i}}\left(\frac{\partial_{0}^{*}}{\partial x_{j}}\left(\frac{1}{r}\right)\right)=\mathcal{P} f\left(\frac{3 x_{i} x_{j}-r^{2} \delta_{i j}}{r^{5}}\right)-4 \pi n_{i} n_{j} \delta_{*} \tag{5.9}
\end{equation*}
$$

The fact that $\frac{\partial_{0}^{*}}{\partial x_{i}}\left(\frac{\partial_{0}^{*}}{\partial x_{j}}\right) \neq \frac{\partial_{0}^{* 2}}{\partial x_{i} \partial x_{j}}$ is obvious in the Example 5 .1. but it is harder to see it in cases like this one. In fact, the fact that the two results are different is overlooked in 10. Observe that the projection of both $4 \pi\left(\delta_{i j}-4 n_{i} n_{j}\right) \delta_{*}$ and of $-4 \pi n_{i} n_{j} \delta_{*}$ onto $\mathcal{D}^{\prime}\left(\mathbb{R}^{3}\right)$ is given by $-(4 \pi / 3) \delta_{i j} \delta(\mathbf{x})$, but this does not mean that they are equal; observe also that one needs the finite part in (5.8) and in (5.9) since the principal value, as used in (5.7), exists in $\mathcal{D}^{\prime}\left(\mathbb{R}^{3}\right)$ but not in $\left(\mathcal{D}_{*}^{[0]}\left(\mathbb{R}^{3}\right)\right)^{\prime}$.

## Acknowledgement

The authors gratefully acknowledge support from NSF, through grant number 0968448.

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Received by the editors September 16, 2013


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